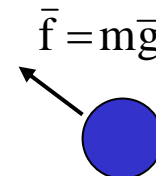
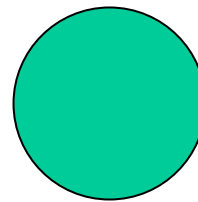
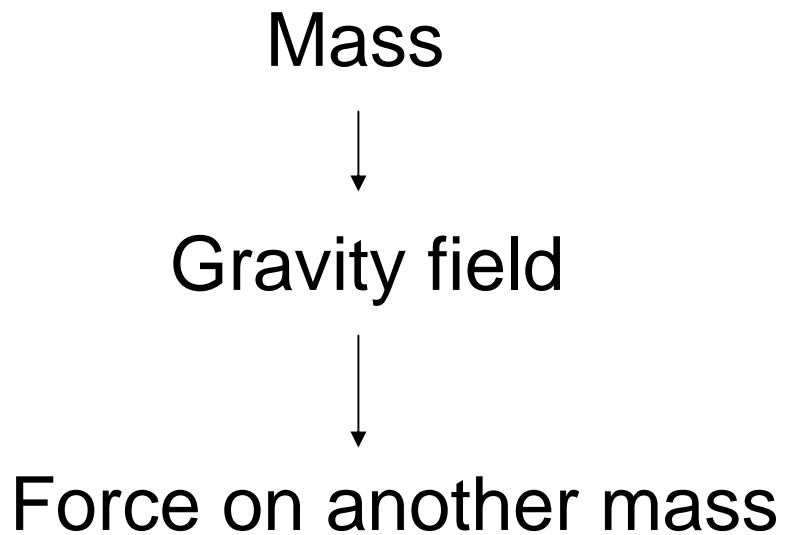


Fundamental electrophysics and engineering design of MAGLEV vehicles

Prof. Jim Hoburg
Department of Electrical & Computer Engineering
Carnegie Mellon University

Gravity



Electricity

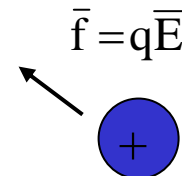
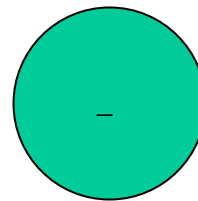
Charge



Electric field



Force on another charge



Magnetism

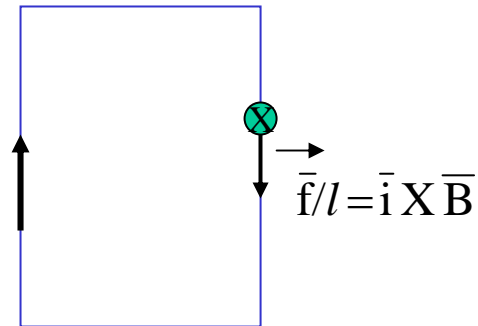
Current



Magnetic field



Force on another
Current segment



Magnetism (microscopic)

Microscopic
current loops

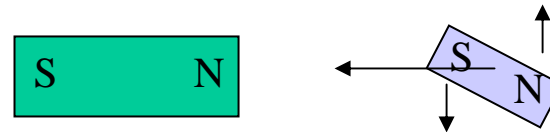


Magnetic field



Force on other

Microscopic
current loops



Macroscopic description of
assembly of microscopic forces:

Net force: opposite poles attract,
same poles repel.

Torque: aligns to imposed field.

Sources \rightarrow Fields

Charge \rightarrow Electric field

$$\rho = \text{div}(\epsilon_0 \bar{E})$$

Current \rightarrow Magnetic field

$$\bar{J} = \text{curl}\left(\frac{\bar{B}}{\mu_0}\right)$$

Sources \rightarrow Fields

(the whole story, including time varying fields as sources of fields)

Maxwell's Equations

$$\rho = \text{div}(\epsilon_0 \bar{E})$$

$$\bar{J} + \epsilon_0 \frac{\partial \bar{E}}{\partial t} = \text{curl}\left(\frac{\bar{B}}{\mu_0}\right)$$

$$0 = \text{div}(\bar{B})$$

$$-\frac{\partial \bar{B}}{\partial t} = \text{curl}(\bar{E})$$

Electromagnetic Waves

Even in a vacuum, where $\rho = 0$ and $\bar{J} = 0$:

$$0 = \text{div}(\epsilon_0 \bar{E})$$

$$0 = \text{div}(\bar{B})$$

$$\epsilon_0 \frac{\partial \bar{E}}{\partial t} = \text{curl} \left(\frac{\bar{B}}{\mu_0} \right)$$

$$-\frac{\partial \bar{B}}{\partial t} = \text{curl}(\bar{E})$$

Electromagnetic Waves

(continued)

$$\begin{array}{l} \downarrow \\ \epsilon_0 \frac{\partial \bar{E}}{\partial t} = \text{curl} \left(\frac{\bar{B}}{\mu_0} \right) \\ \downarrow \\ -\frac{\partial \bar{B}}{\partial t} = \text{curl}(\bar{E}) \end{array}$$

$$\epsilon_0 = 8.854 \times 10^{-12} \frac{\text{Farad}}{\text{m}} \quad (\rho \rightarrow \bar{E})$$

$$\mu_0 = 1.257 \times 10^{-6} \frac{\text{Henry}}{\text{m}} \quad (\bar{J} \rightarrow \bar{B})$$

$$\frac{1}{\sqrt{\mu_0 \epsilon_0}} = 3.00 \times 10^8 \text{ m/s} = c = \text{speed of light in vacuum}$$

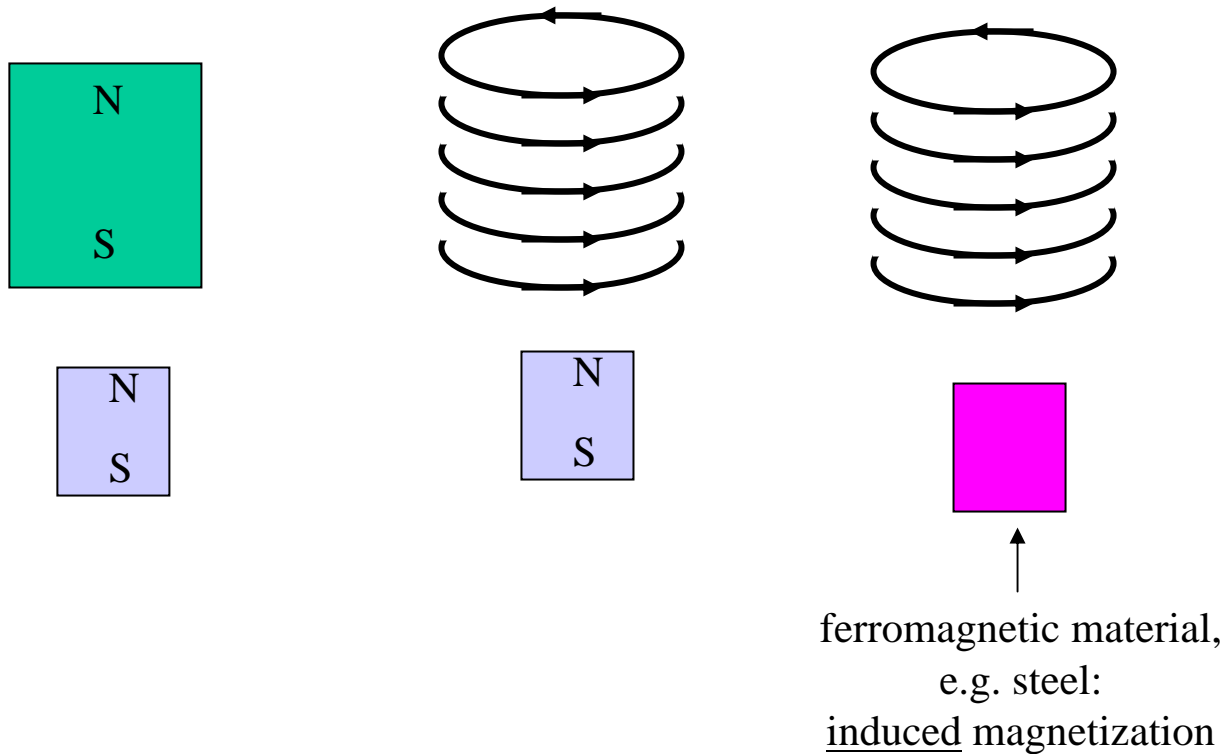
What does this have to do with MAGLEV?

Same fundamental physics describes:

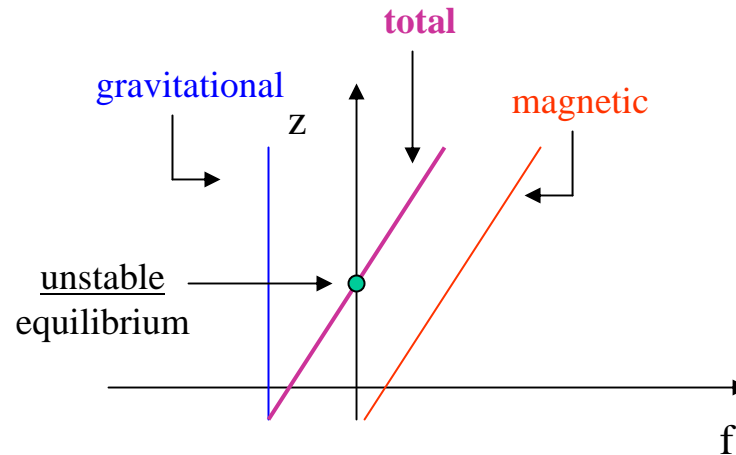
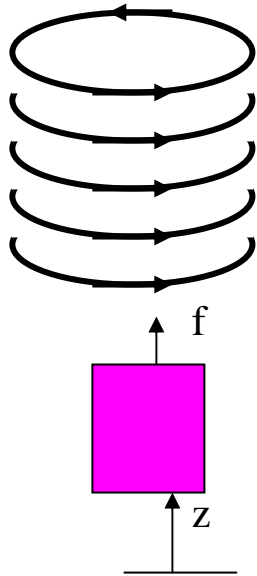
- Light, lasers, X-rays, ... (electromagnetic spectrum, why is the sky blue?)
- Wireless communications (radio, TV, cell phones, wireless computers, ...)
- Integrated circuits (computer chips)
- Lightning
- Electrostatic precipitation
- Electrophotography & laser printers
- Microelectromechanical systems (MEMS)
- Magnetic memory (tapes, disks, MRAM, magnetic stripes, ...)
- Rotating electrical machinery (generators & motors)
- Linear synchronous motor (LSM)
- Magnetic confinement for nuclear fusion
- MAGLEV

Attractive magnetic levitation:

“Electromagnetic levitation”

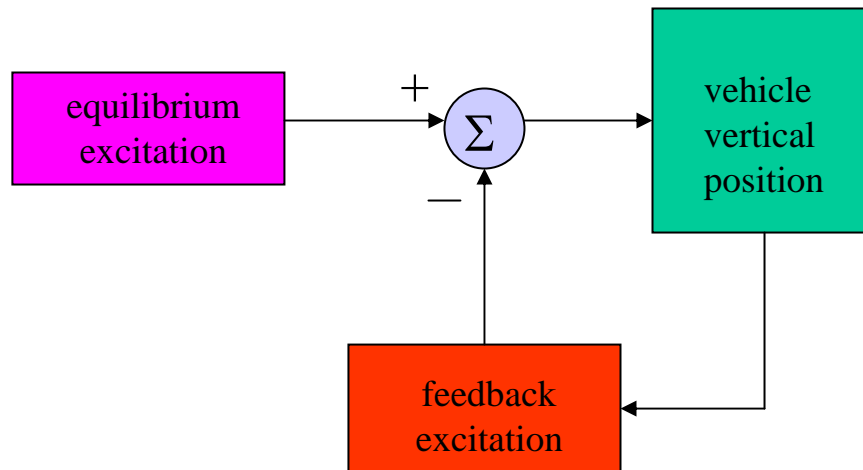


Attractive magnetic levitation:



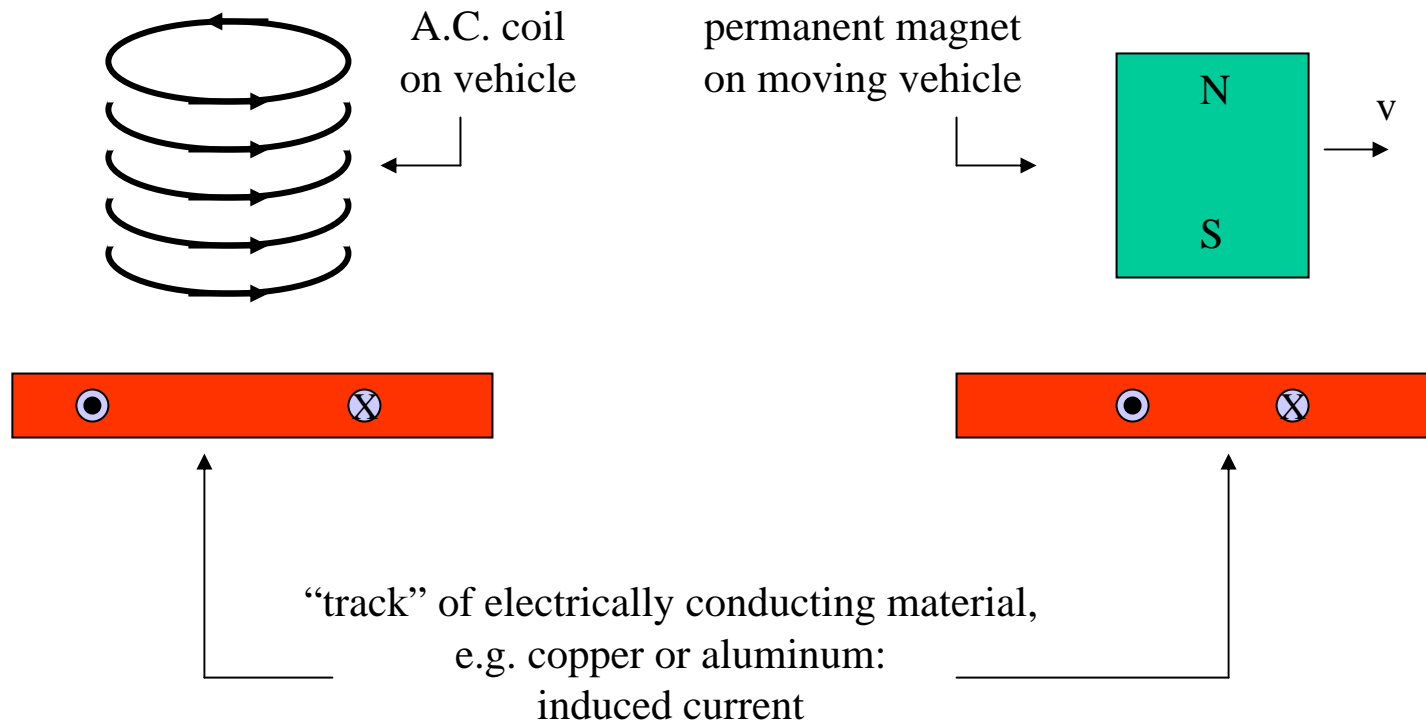
(inherently unstable)

Attractive magnetic levitation: requires feedback control system to stabilize



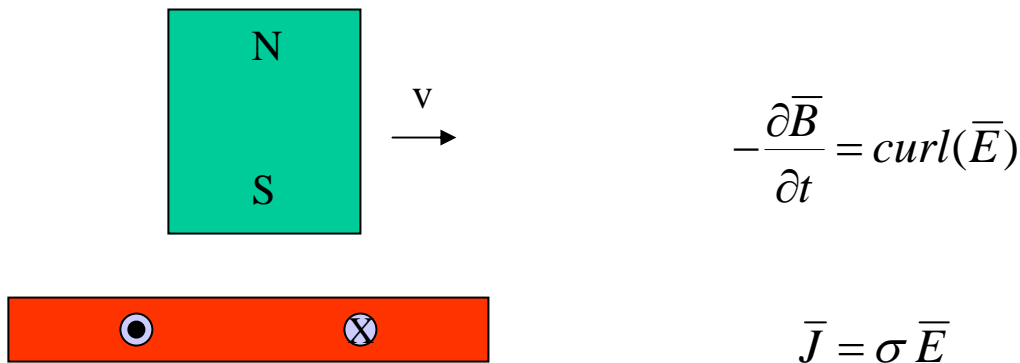
Repulsive magnetic levitation:

“Electrodynamic levitation”

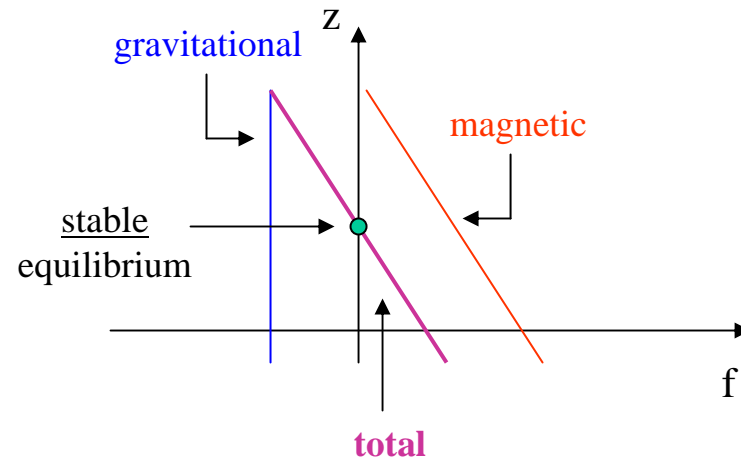
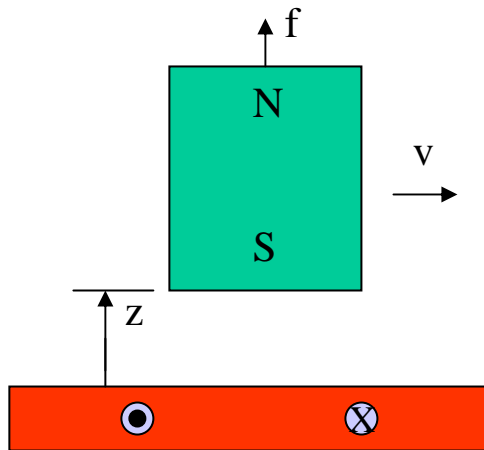


Repulsive magnetic levitation:

This interaction involves an induced electric field in the conducting material, caused by a time-varying imposed magnetic field.



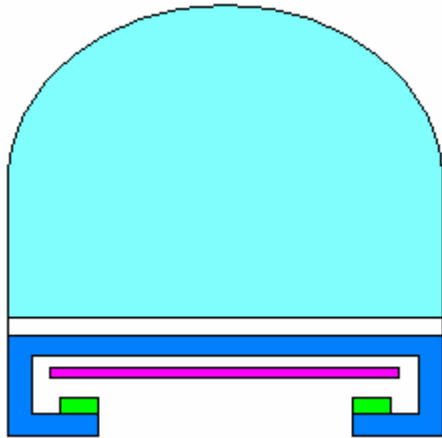
Repulsive magnetic levitation:



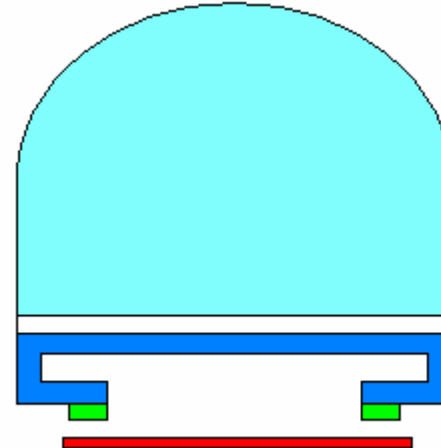
(inherently stable:
no feedback control system needed)

Either mechanism can be used to levitate a vehicle

- Attractive levitation



- Repulsive levitation



High speed (~ 450 km/hr) MAGLEV systems:

- German Transrapid:
 - Attractive (“electromagnetic”) levitation via conventional electromagnets
 - LSM propulsion
- Japanese MLX:
 - Repulsive (“electrodynamic”) levitation via superconducting magnets (on-board cryogenics)
 - LSM propulsion

Transrapid Test Vehicle TR-08



Germany, 1999

Null Flux Suspension Vehicle MLX01

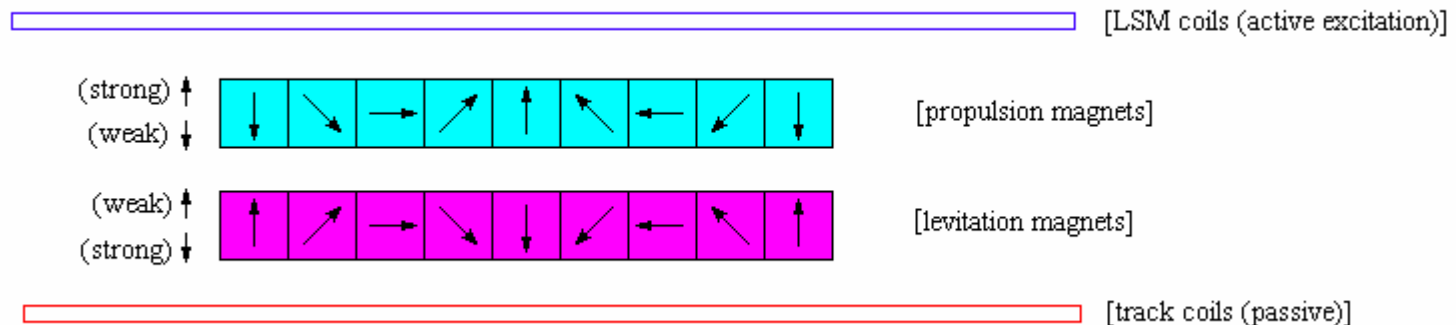


Japan, 1997 (Yamanashi test facility)

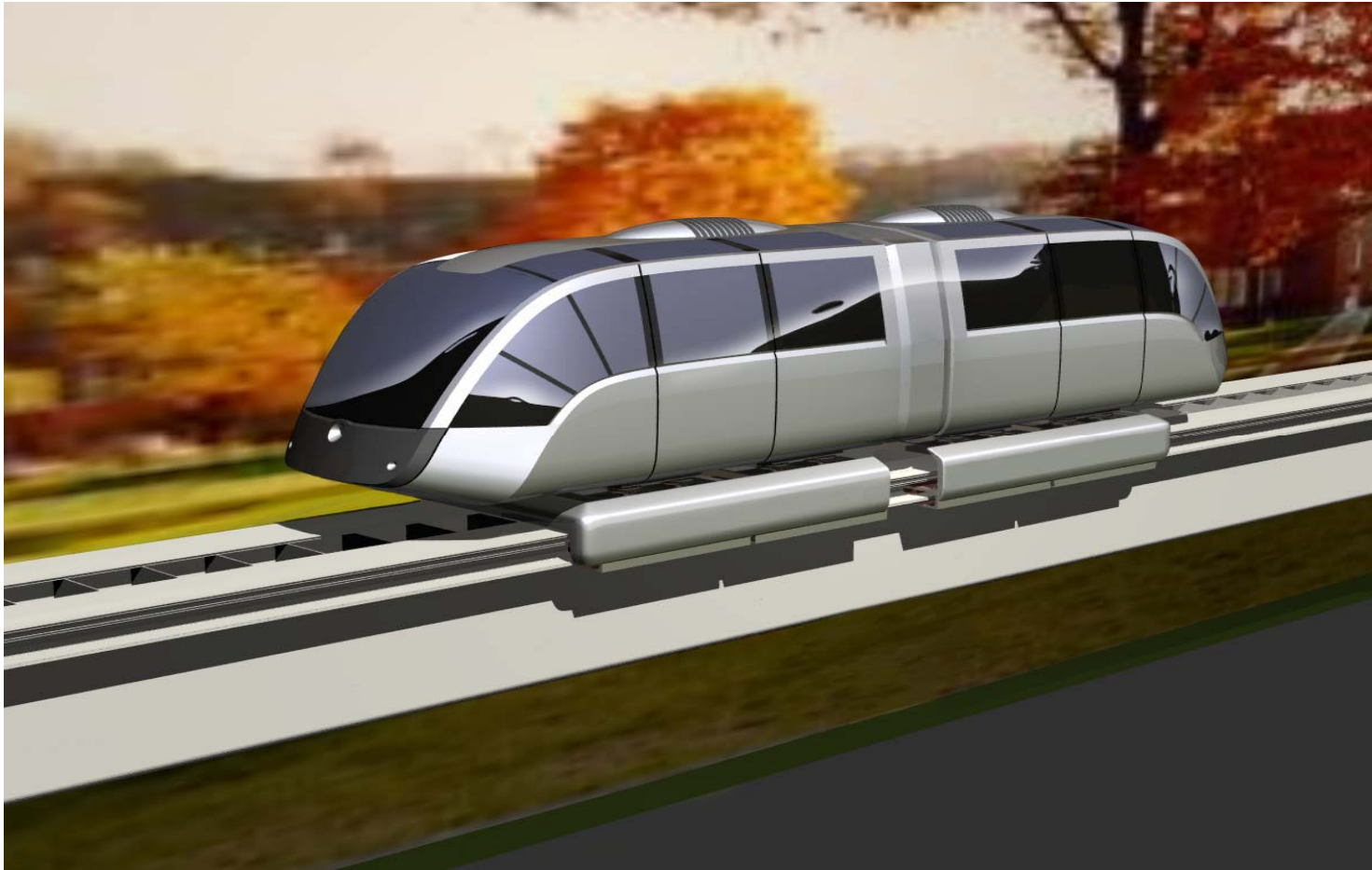
Low speed urban MAGLEV concept:

- Permanent (NdFeB) magnet arrays on vehicle:
 - Repulsive (“electrodynamic”) levitation via induced currents in track coils
 - LSM propulsion

Halbach permanent magnet arrays



(Envisioned) Pittsburgh Urban Maglev vehicle

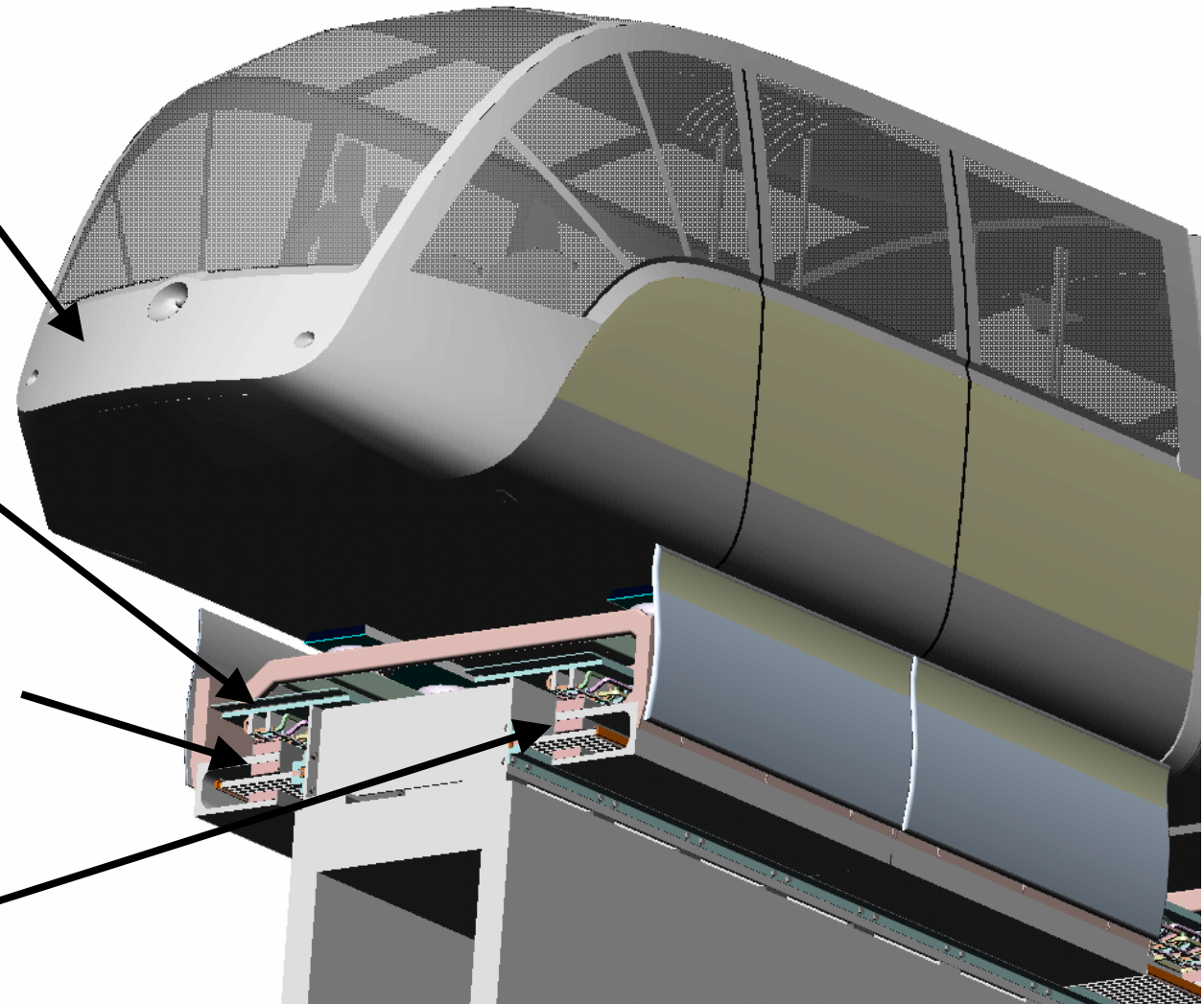


**Vehicle on
Guideway**

**Linear
Synchronous
Motor**

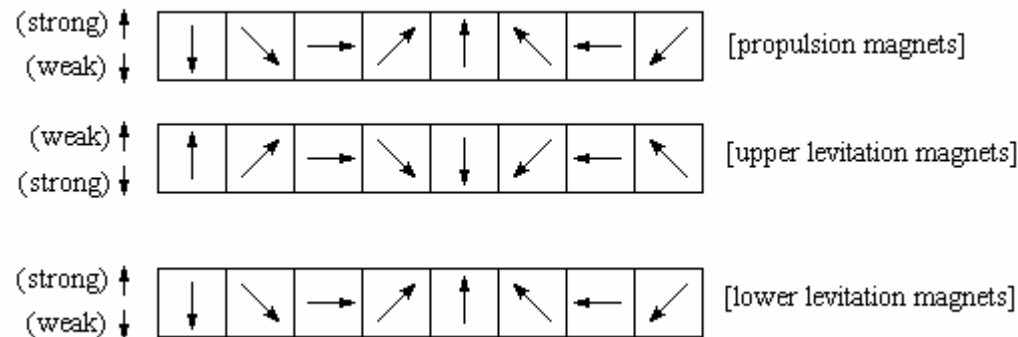
Suspension Track

**Double Sided
Magnet Array**



Magnetic fields from known sources:

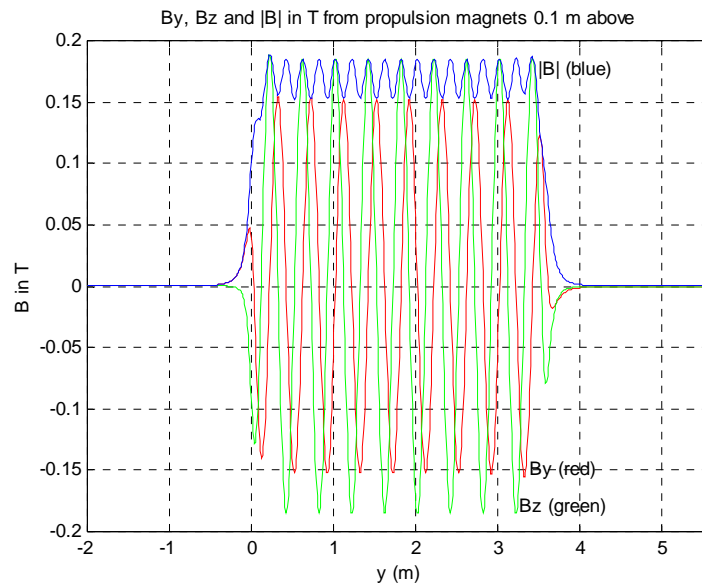
computing passenger compartment field levels



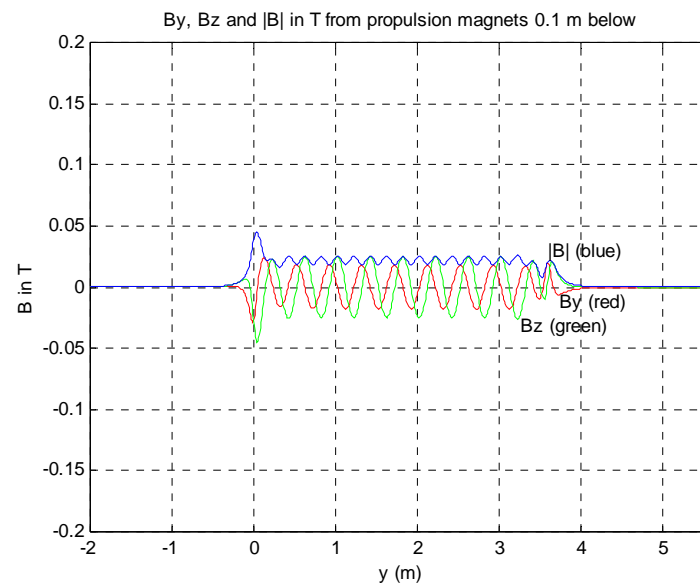
Halbach array fields: basic structure via magnetization charge description:

- propulsion magnet near fields (0.1 m above & below):

components & magnitude above



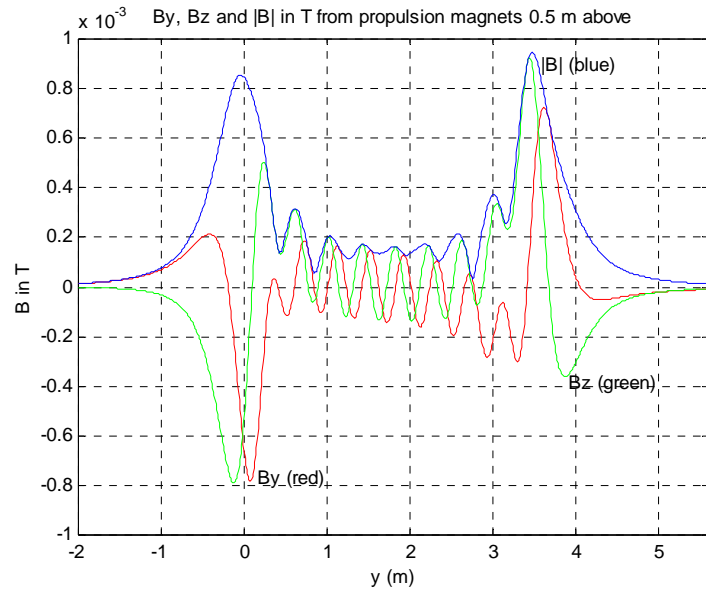
components & magnitude below



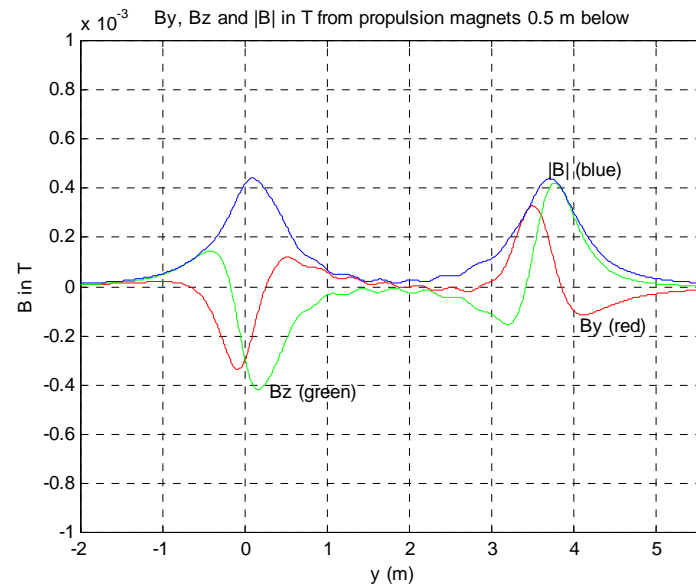
Halbach array fields: basic structure via magnetization charge description:

- propulsion magnet far fields (0.5 m above & below):

components & magnitude above



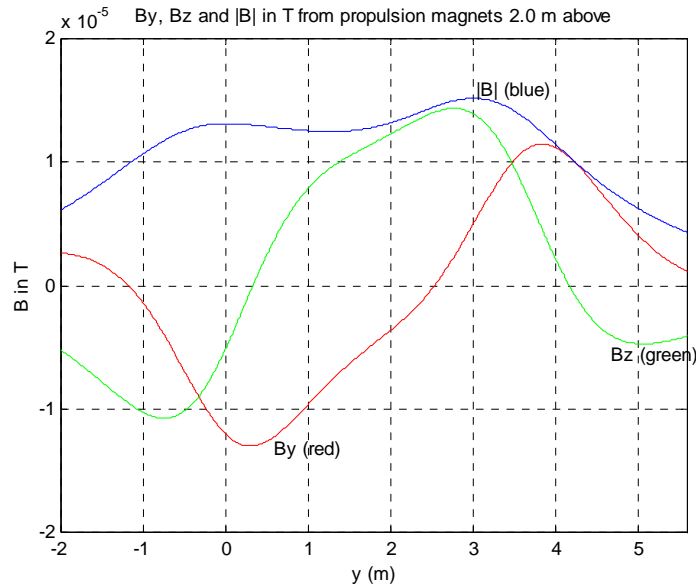
components & magnitude below



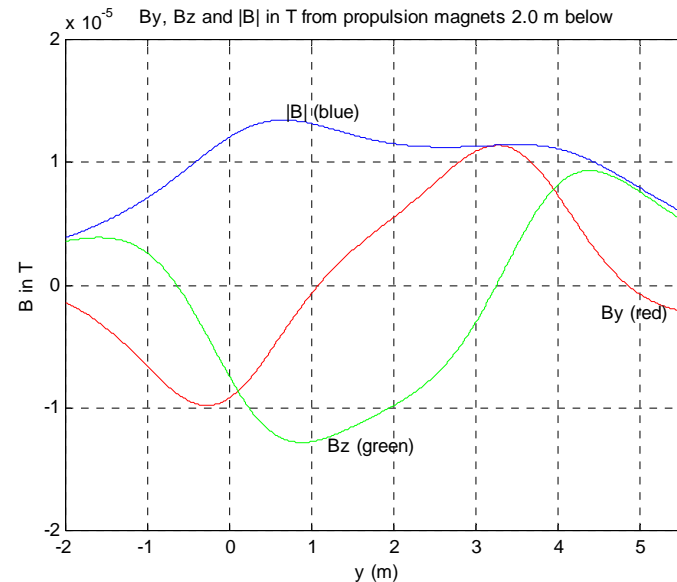
Halbach array fields: basic structure via magnetization charge description:

- propulsion magnet very far fields (2.0 m above & below):

components & magnitude above

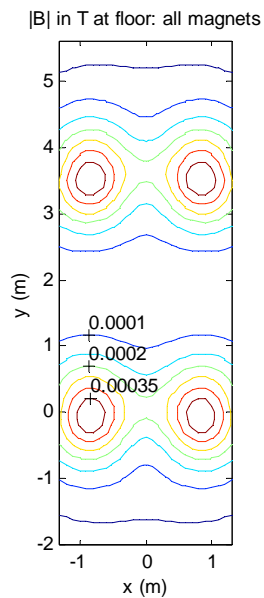


components & magnitude below

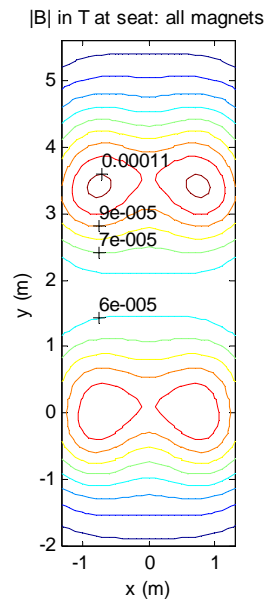


Fields in passenger compartment magnetization charge description:

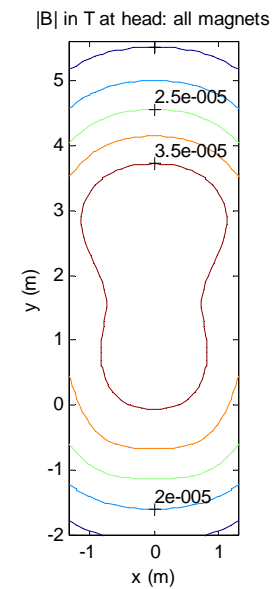
floor level contour plot



seat level contour plot

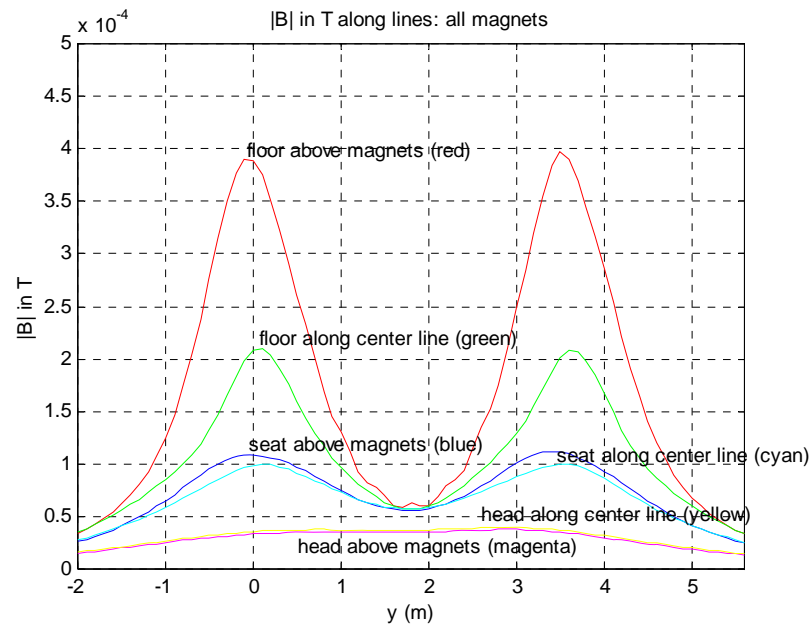


head level contour plot

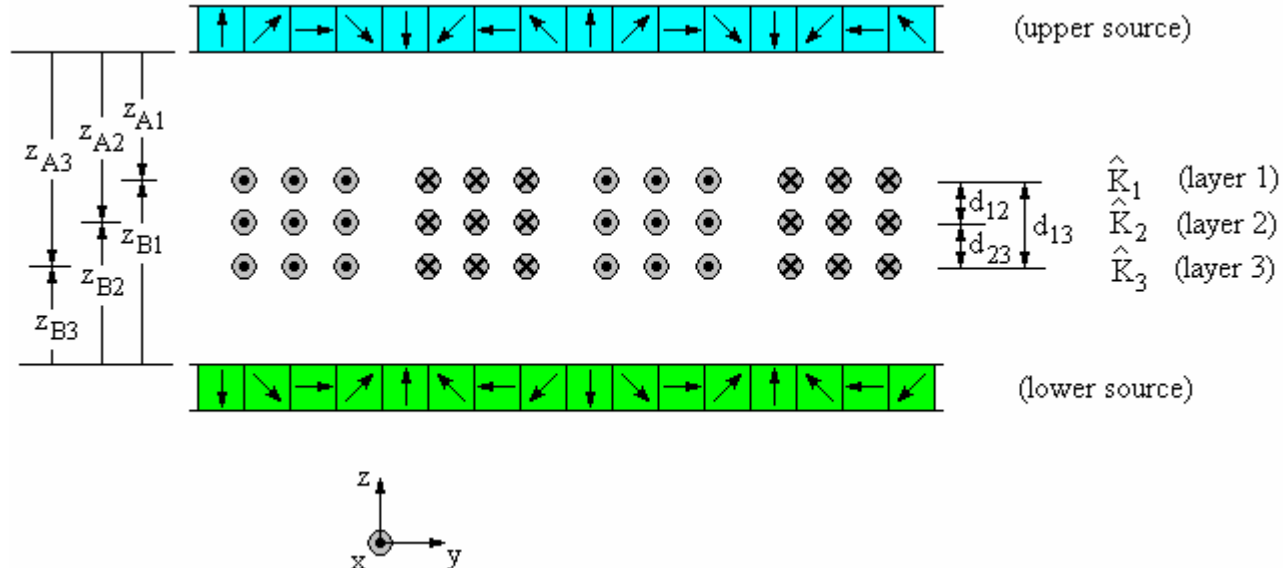


Fields in passenger compartment magnetization charge description:

line plots at floor, seat and head levels



Electromechanics: computing lift & drag versus velocity



Faraday's Law for rectangular contours in lamination planes:

$$\oint_C \bar{E} \cdot d\bar{l} = -\frac{d}{dt} \int_S \mu_0 \bar{H} \cdot \bar{n} da$$

$$2\ell \hat{E}_x = -j\omega \left[\hat{\Lambda}_A + \hat{\Lambda}_B + p \left(\hat{\Lambda}_{self} + \sum \hat{\Lambda}_b + \sum \hat{\Lambda}_a \right) \right]$$

Electric field related to surface currents in lamination planes:

$$E_x = \frac{K_x}{\sigma \Delta}$$

Time constant for induced currents:

$$\tau_m = \frac{\mu_0 \sigma \Delta \lambda}{4\pi}$$

Driving frequency based upon vehicle velocity:

$$\omega = 2\pi \frac{v}{\lambda}$$

All induced currents in laminations take the forms:

$$\bar{K} = \bar{i}_x \operatorname{Re} \left\{ \hat{K} \exp \left[j \frac{2\pi}{\lambda} (vt - y) \right] \right\}$$

Resultant fields (above and below) any one lamination are:

$$\hat{H}_a(y, z) = \frac{\hat{K}}{2} [-\bar{i}_y + (j)\bar{i}_z] e^{-jky} e^{-kz}$$

$$\hat{H}_b(y, z) = \frac{\hat{K}}{2} [+ \bar{i}_y + (j)\bar{i}_z] e^{-jky} e^{+kz}$$

Simultaneous equations governing induced currents in laminations:

$$\begin{bmatrix} \left(1 + \frac{1}{j\omega p \tau_m}\right) & e^{-kd_{12}} & e^{-kd_{13}} & \dots \\ e^{-kd_{12}} & \left(1 + \frac{1}{j\omega p \tau_m}\right) & e^{-kd_{23}} & \dots \\ e^{-kd_{13}} & e^{-kd_{23}} & \left(1 + \frac{1}{j\omega p \tau_m}\right) & \dots \\ \dots & \dots & \dots & \dots \end{bmatrix} \begin{bmatrix} \hat{K}_1 \\ \hat{K}_2 \\ \hat{K}_3 \\ \dots \end{bmatrix}$$

$$= \frac{j2}{p\ell} \begin{bmatrix} \int_{-\ell/2}^{\ell/2} \hat{H}_z^A(x, z_{A1}) dx + \int_{-\ell/2}^{\ell/2} \hat{H}_z^B(x, z_{B1}) dx \\ \int_{-\ell/2}^{\ell/2} \hat{H}_z^A(x, z_{A2}) dx + \int_{-\ell/2}^{\ell/2} \hat{H}_z^B(x, z_{B2}) dx \\ \int_{-\ell/2}^{\ell/2} \hat{H}_z^A(x, z_{A3}) dx + \int_{-\ell/2}^{\ell/2} \hat{H}_z^B(x, z_{B3}) dx \\ \dots \end{bmatrix} \equiv \frac{j2}{p\ell} \begin{bmatrix} \hat{I}_z^1 \\ \hat{I}_z^2 \\ \hat{I}_z^3 \\ \dots \end{bmatrix}$$

Time-averaged lift and drag forces (per wavelength) on vehicle:

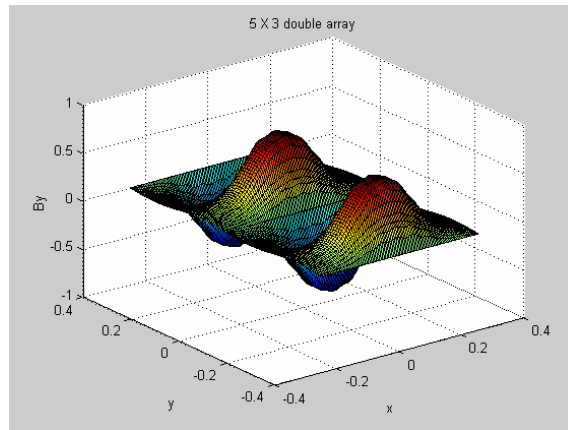
$$L_\lambda = -\frac{\mu_0 \lambda p}{2} \sum_i \operatorname{Re} \left\{ \int_{-\ell/2}^{\ell/2} \hat{K}_i^* \left[\hat{H}_y^A(x, z_{Ai}) + \hat{H}_y^B(x, z_{Bi}) \right] dx \right\} \equiv -\frac{\mu_0 \lambda p}{2} \sum_i \operatorname{Re} \left\{ \hat{K}_i^* \hat{I}_y^i \right\}$$

$$D_\lambda = -\frac{\mu_0 \lambda p}{2} \sum_i \operatorname{Re} \left\{ \int_{-\ell/2}^{\ell/2} \hat{K}_i^* \left[\hat{H}_z^A(x, z_{Ai}) + \hat{H}_z^B(x, z_{Bi}) \right] dx \right\} \equiv -\frac{\mu_0 \lambda p}{2} \sum_i \operatorname{Re} \left\{ \hat{K}_i^* \hat{I}_z^i \right\}$$

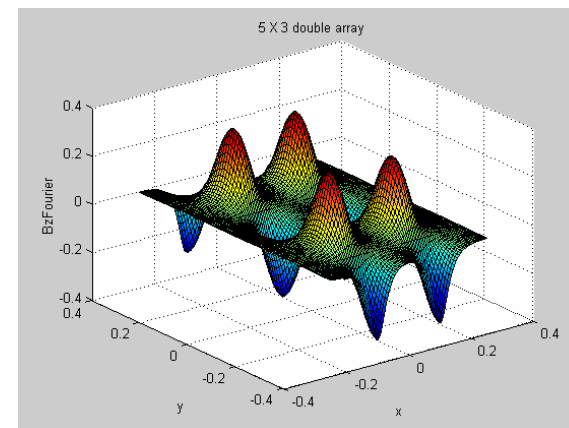
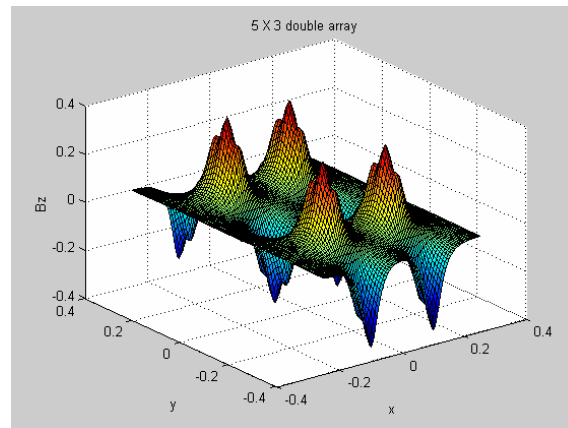
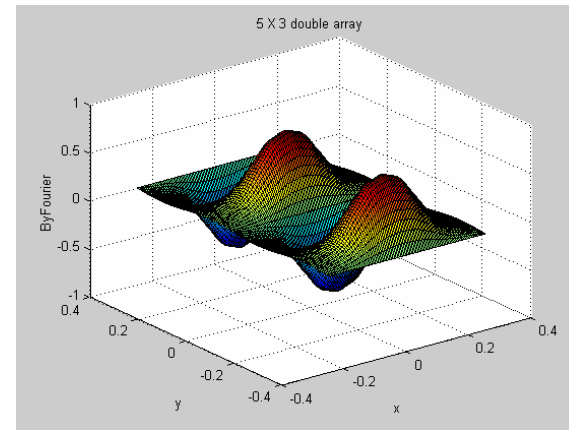
Horizontal (y) and vertical (z) field components are based upon 3-D magnetization charge description.

Complex amplitudes are based upon first Fourier components of y dependences at each value of x.

Double (5 above X 3 below) array
Full 3-D magnetization charge fields:

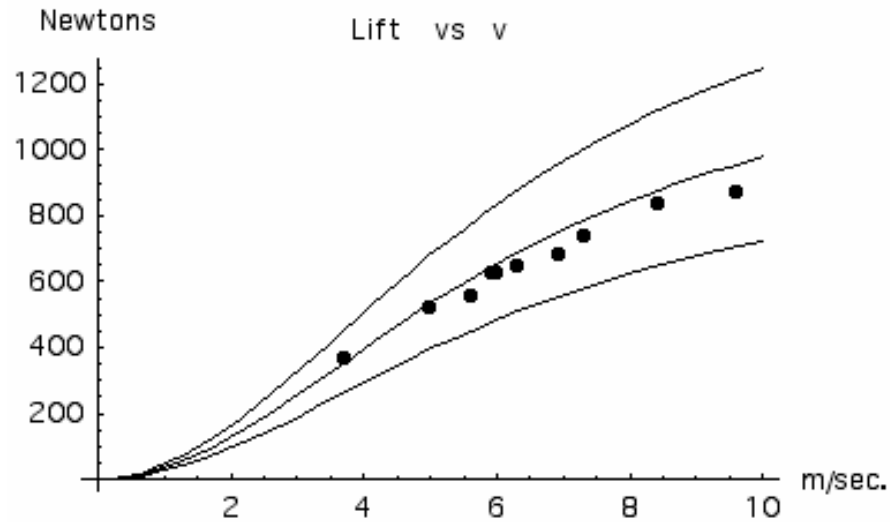


Double (5 above X 3 below) array
First Fourier component approximations:

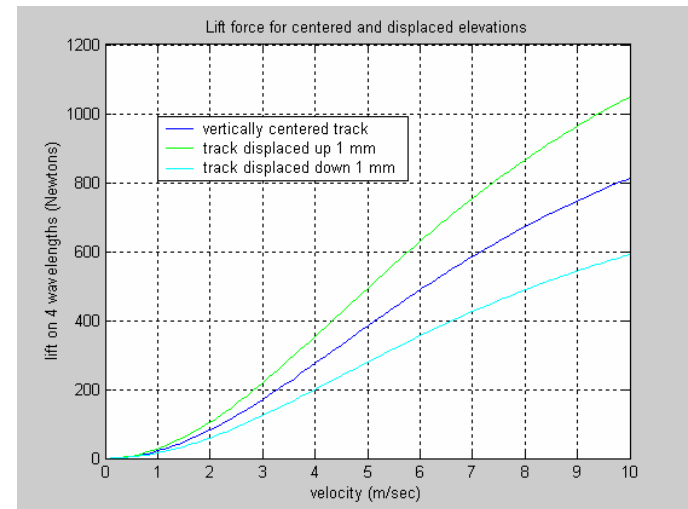


Comparison with Dick Post's model for LLNL test rig

Post:

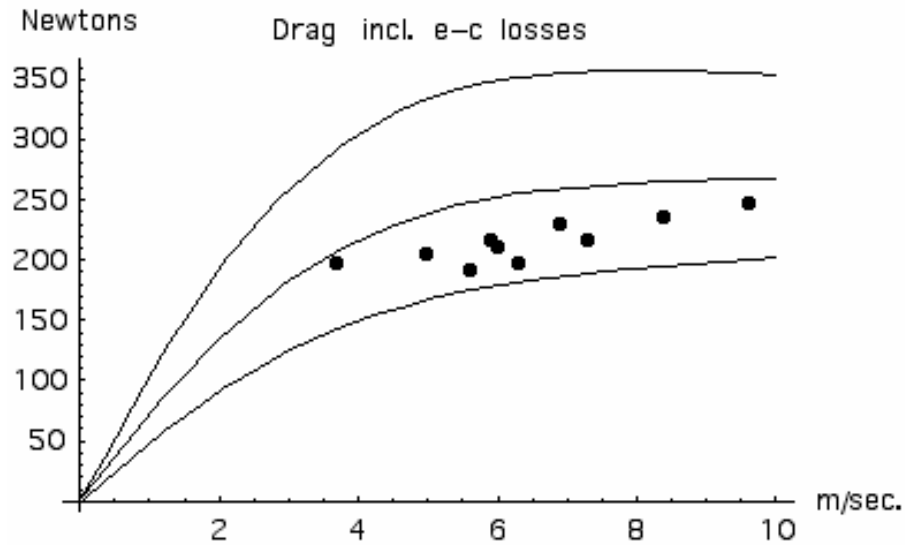


Hoburg:



Comparison with Dick Post's model for LLNL test rig

Post:



Hoburg:

