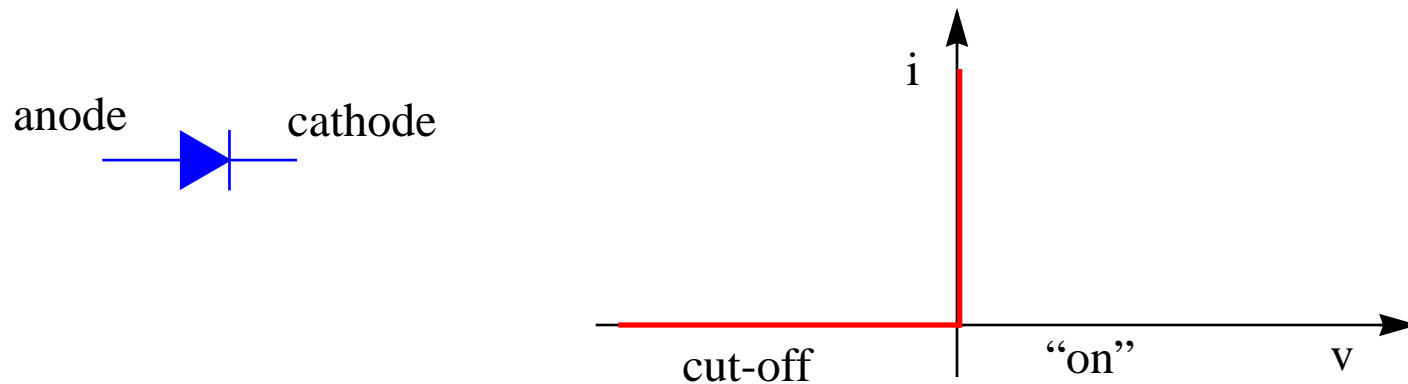


# Diodes

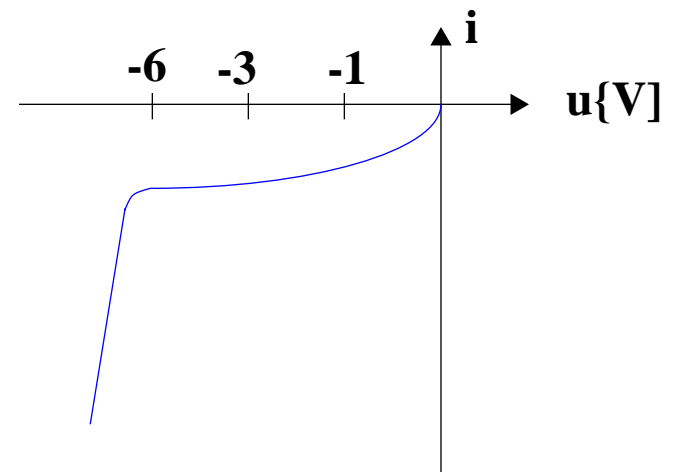
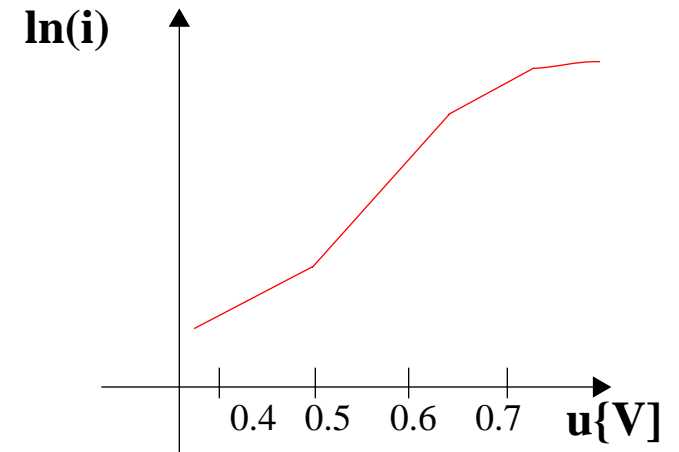
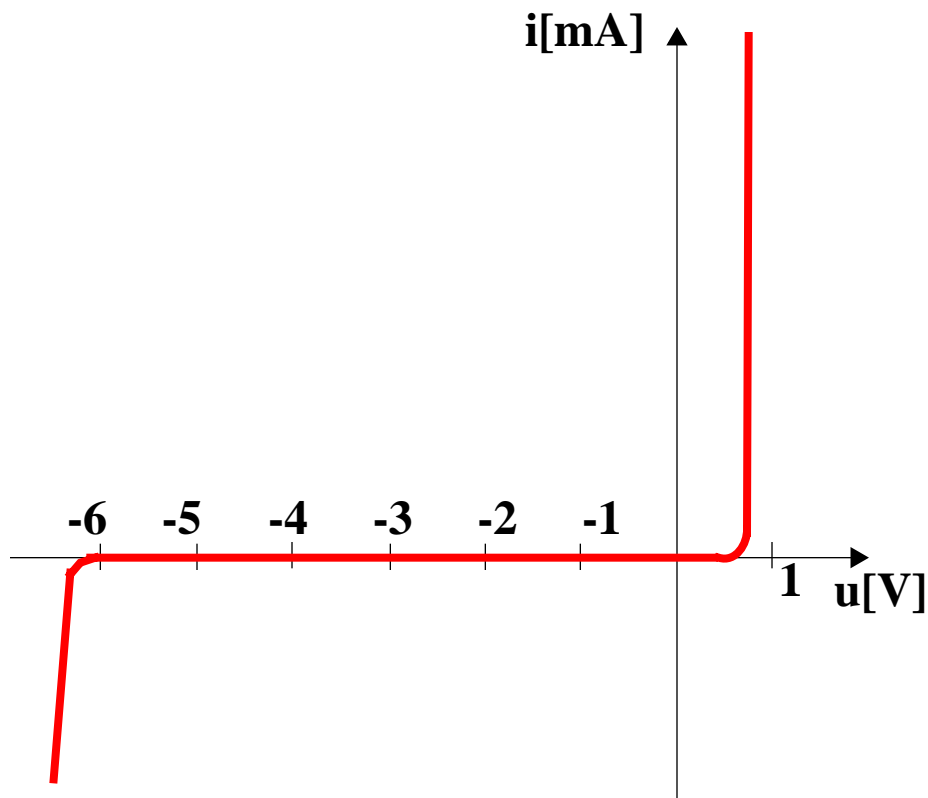
- Simplest nonlinear circuit element
- Basic operation sets the foundation for **Bipolar Junction Transistors (BJTs)**
- Also present in **Field Effect Transistors (FETs)**
- Ideal diode characteristic



- Can be approximated by a **piecewise-linear-like** characteristic

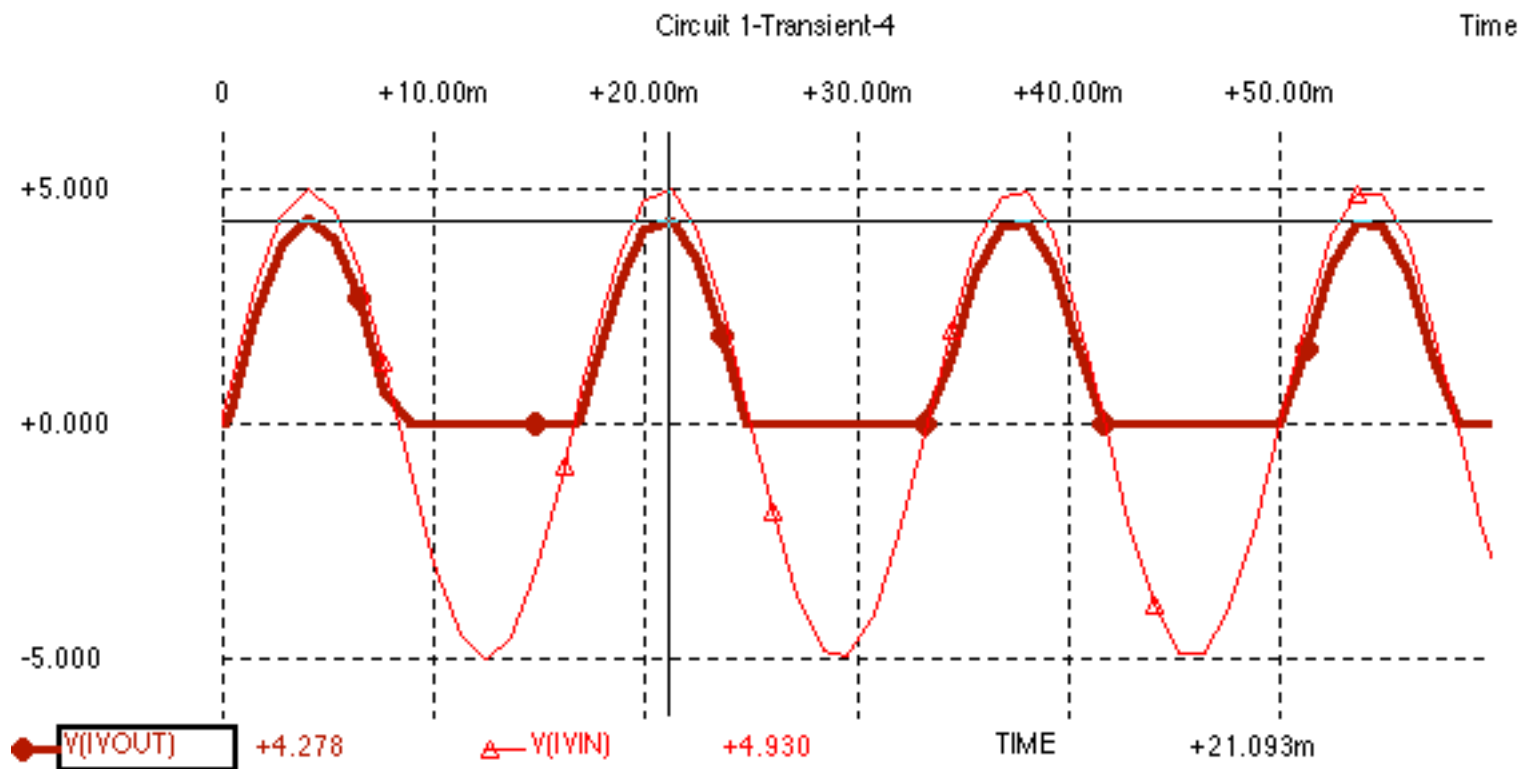
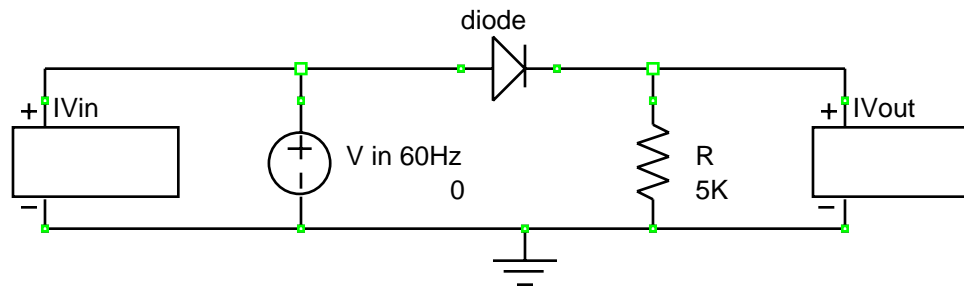
## Actual p-n Diode Characteristic

- An actual diode does not follow this ideal behavior
- The “turn-on” voltage is generally between 0.5 and 0.8 volts
- The diode **breaks down** for large reverse bias



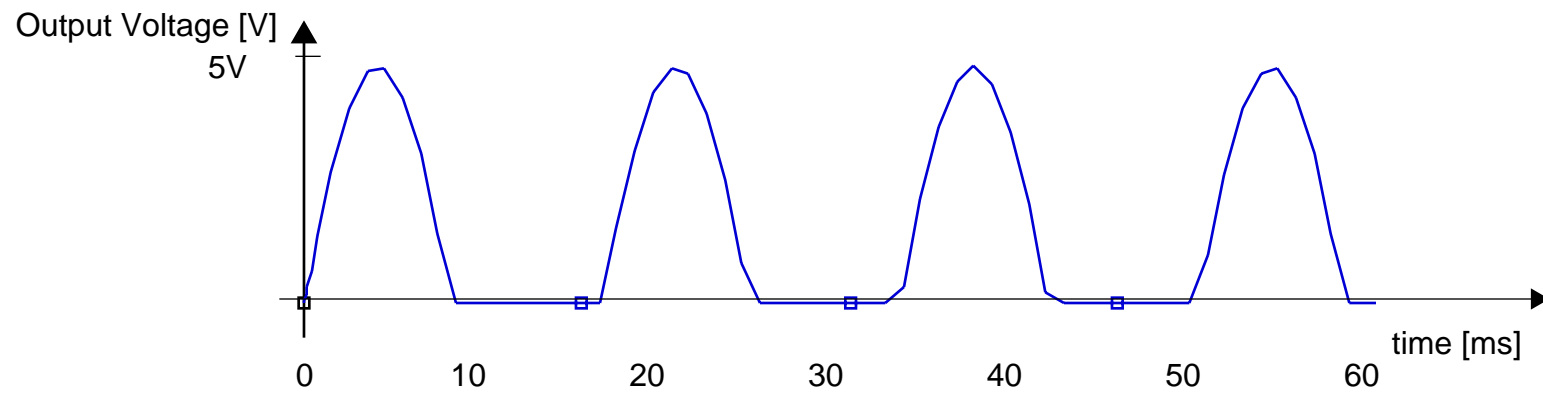
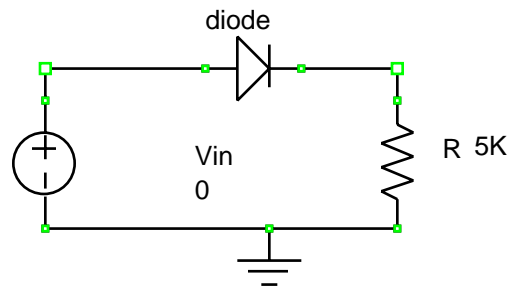
# Diodes for Rectification

- Commonly used for power supplies and to convert ac signals to dc signals



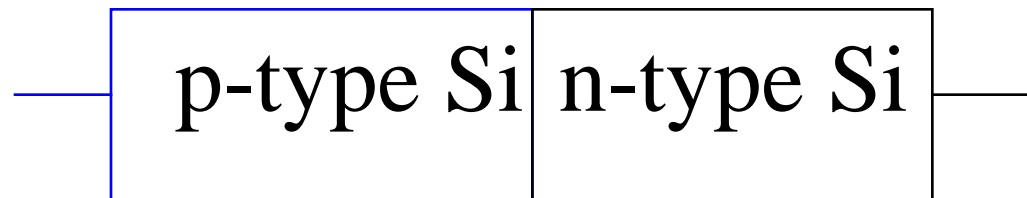
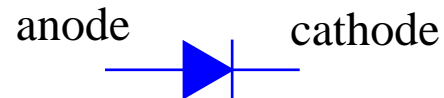
# Diodes for Rectification

- This is a **half-wave** recitifier



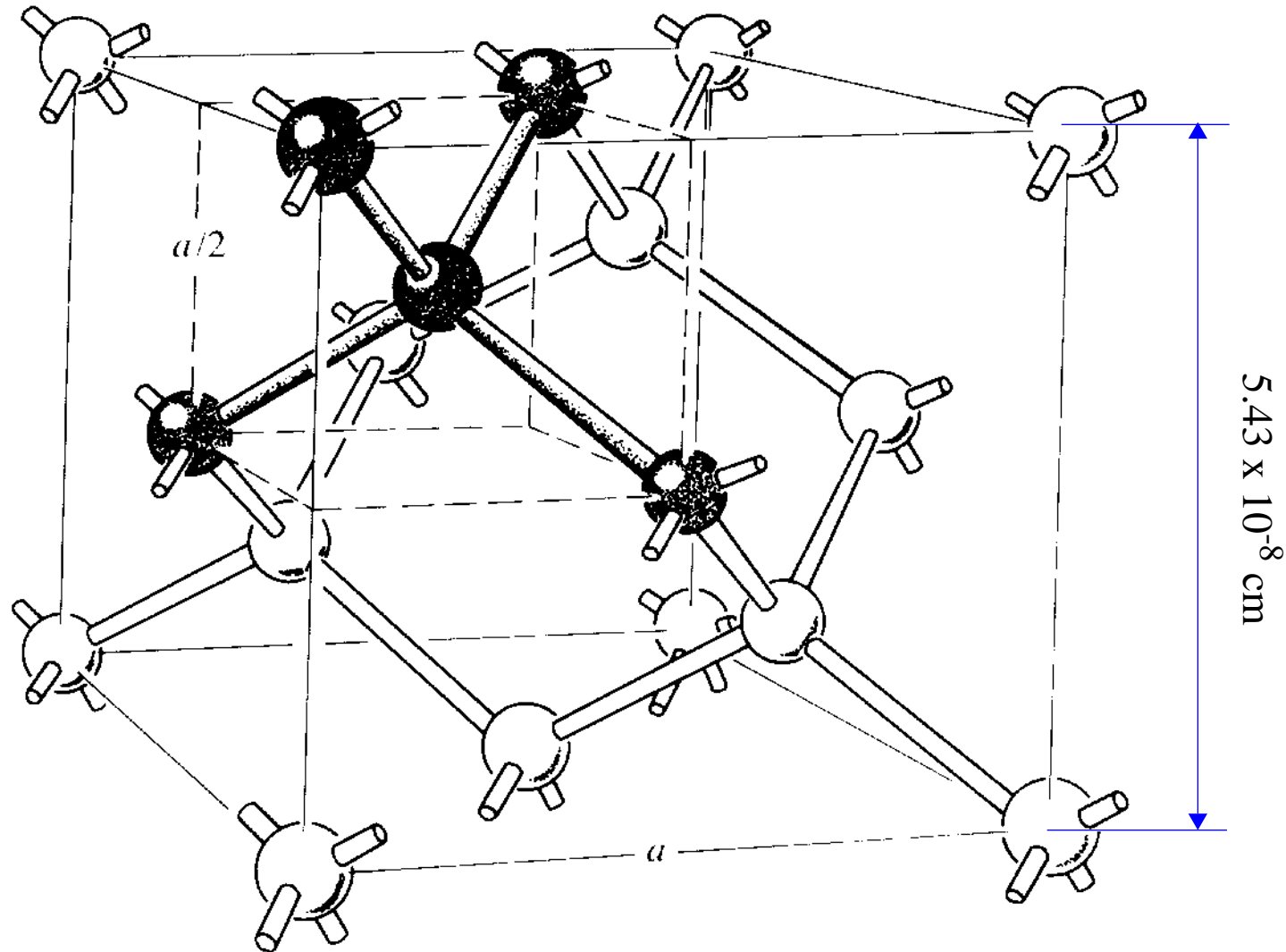
# Silicon Diodes

- Different types of diodes, but we're interested in semiconductor silicon diodes
- Si diodes are formed by the junction of n- and p-type silicon



- What is a semiconductor?
  - “a material which can have electrons and holes as charge carriers”  
(naive definition, but useful for our purpose)
- What are holes?
- What is n-type and p-type silicon?

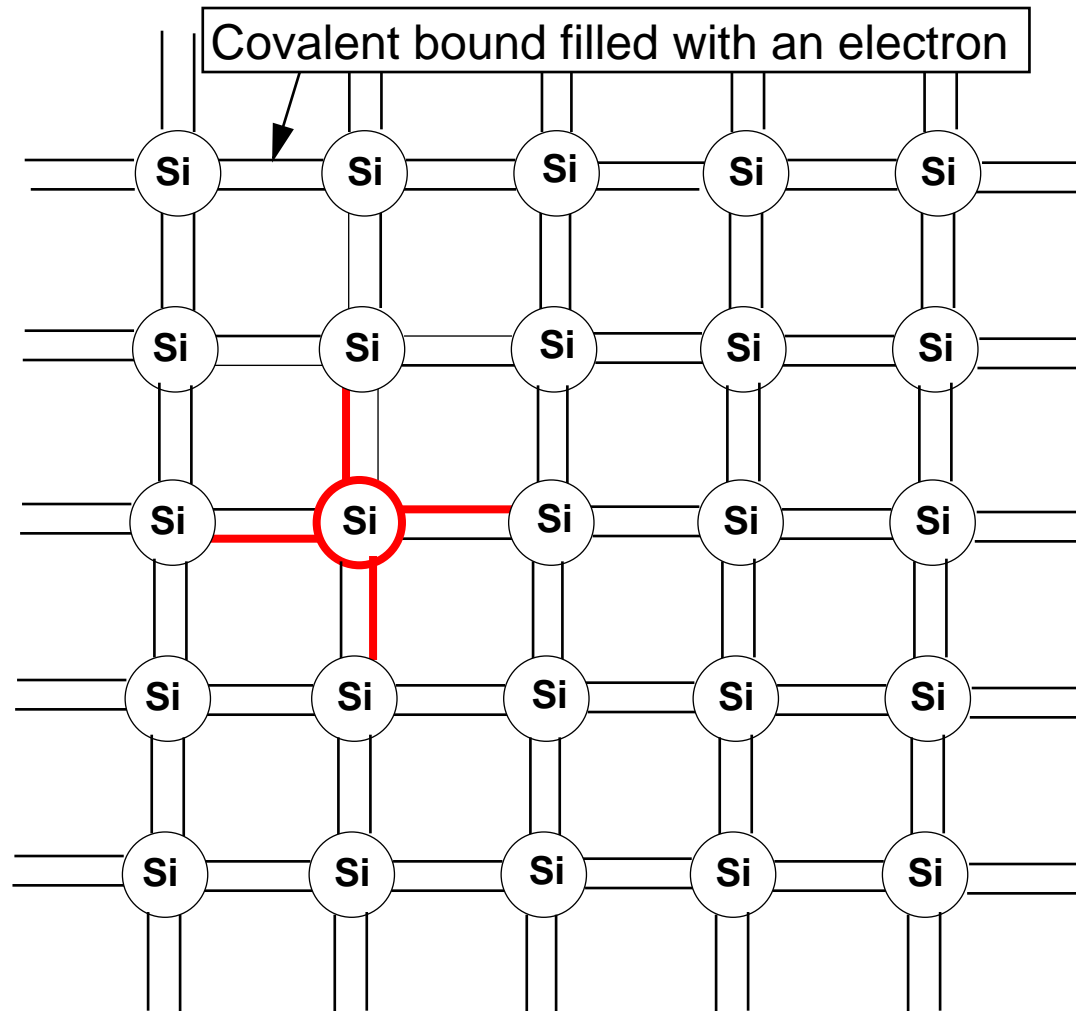
## Electrons and Holes(1) - Mono crystal of Silicon



- In this regular lattice, each Si atom is bound with its four nearest neighbors.
- This ideal crystal lattice looks the same from each node.

## Electrons and Holes(2) -

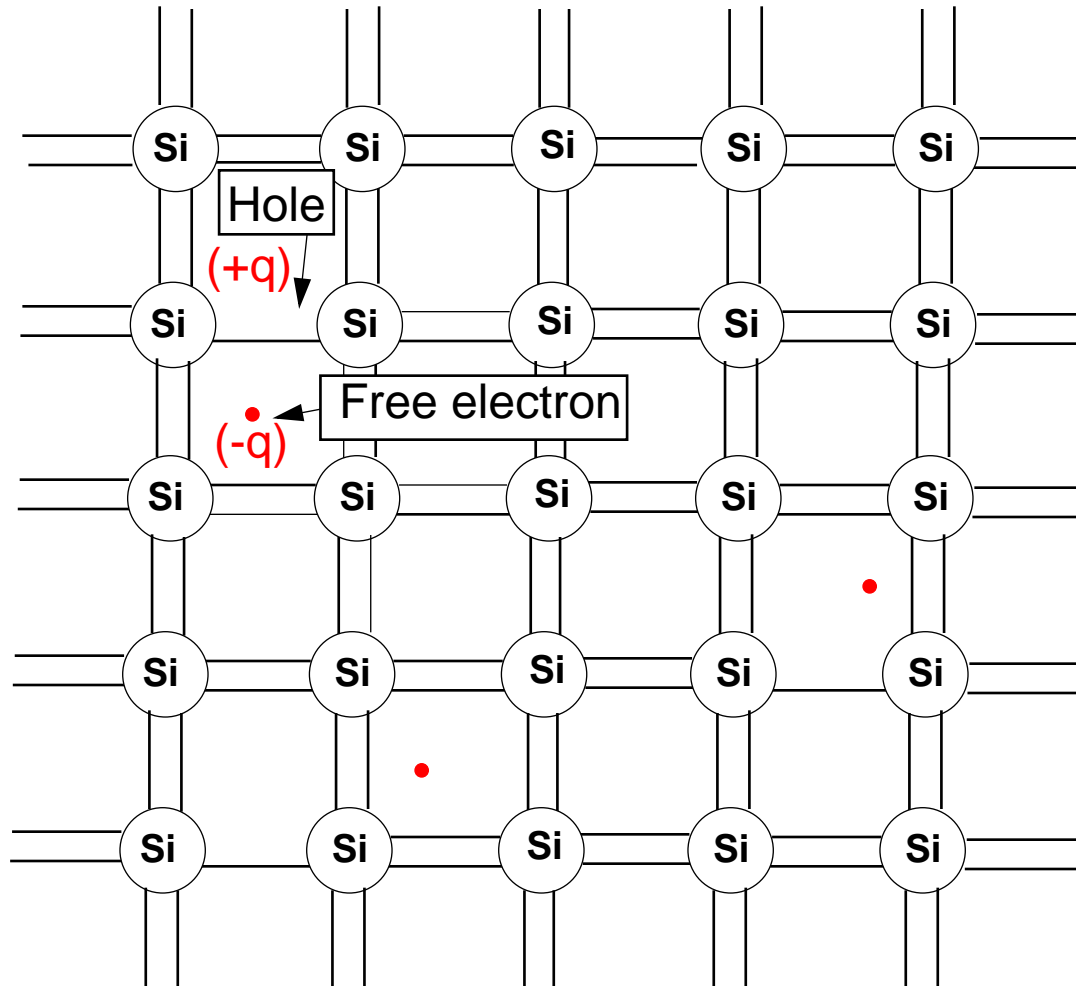
## Two-dimensional silicon lattice



There is about  $5 \times 10^{22}$  Si atoms in  $1 \text{ cm}^3$

## Electrons and Holes(3) -

## Generation and recombination



- Free electrons and holes:
  - appear (**generation**)
  - disappear (**recombination**)

- The number of electron-hole pairs **increases with temperature**. It depends upon the number of electrons which have energy high enough to break free from the bounds

In pure Si crystal at room temperature (**300 K**), there is about  $1.4 \times 10^{10}$  electron-hole pairs in  $1 \text{ cm}^3$ . This is so-called intrinsic concentration  $n_i$ .

In the absolute zero (**0 K**) there would be no such pairs at all.



## Electrons and Holes(4) - Drift Current

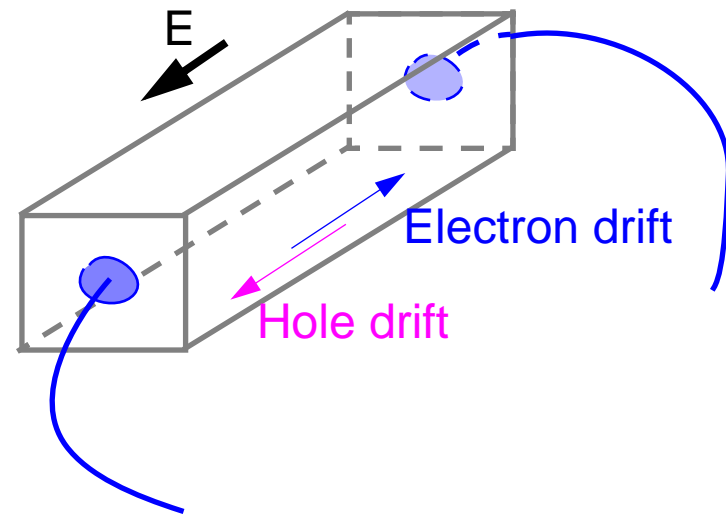
- A voltage applied to a silicon sample produces an electric field that causes the free electrons and holes to **drift** with an average velocity (different for electrons than for holes).
- Assuming we know the average velocity of the electrons in the electric field, we can easily calculate the current density  $J$  (current per unit area):  
 $J = (\text{free charge concentration}) * (\text{drift velocity})$

$v = \mu E$  - average drift velocity

**Holes:**  $J_p (\text{drift}) = q\mu_p p E$

**Electrons:**  $J_n (\text{drift}) = q\mu_n n E$

$J = J_n + J_p$   $J = q(\mu_n n + \mu_p p) E$

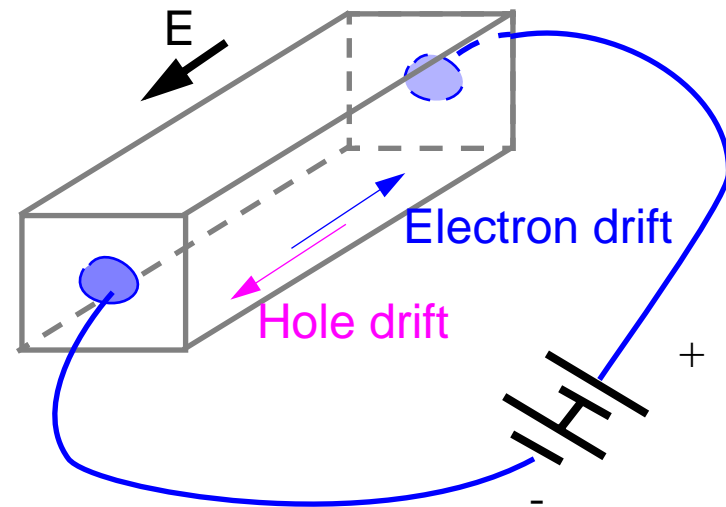


- $p, n$  - concentration of holes and electrons [ $\text{cm}^{-3}$ ]
- $\mu_p, \mu_n$  - mobility of holes and electrons [ $\text{cm}^2/\text{Vs}$ ]
- $q$  - electronic charge (magnitude) =  $1.6 \times 10^{-19}$  [Q]

## Electrons and Holes(5) - Drift Current

- What is the current through this sample then?

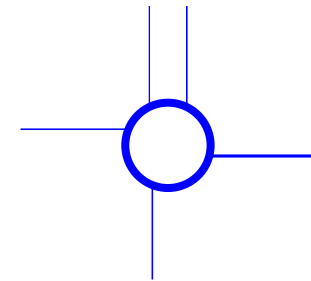
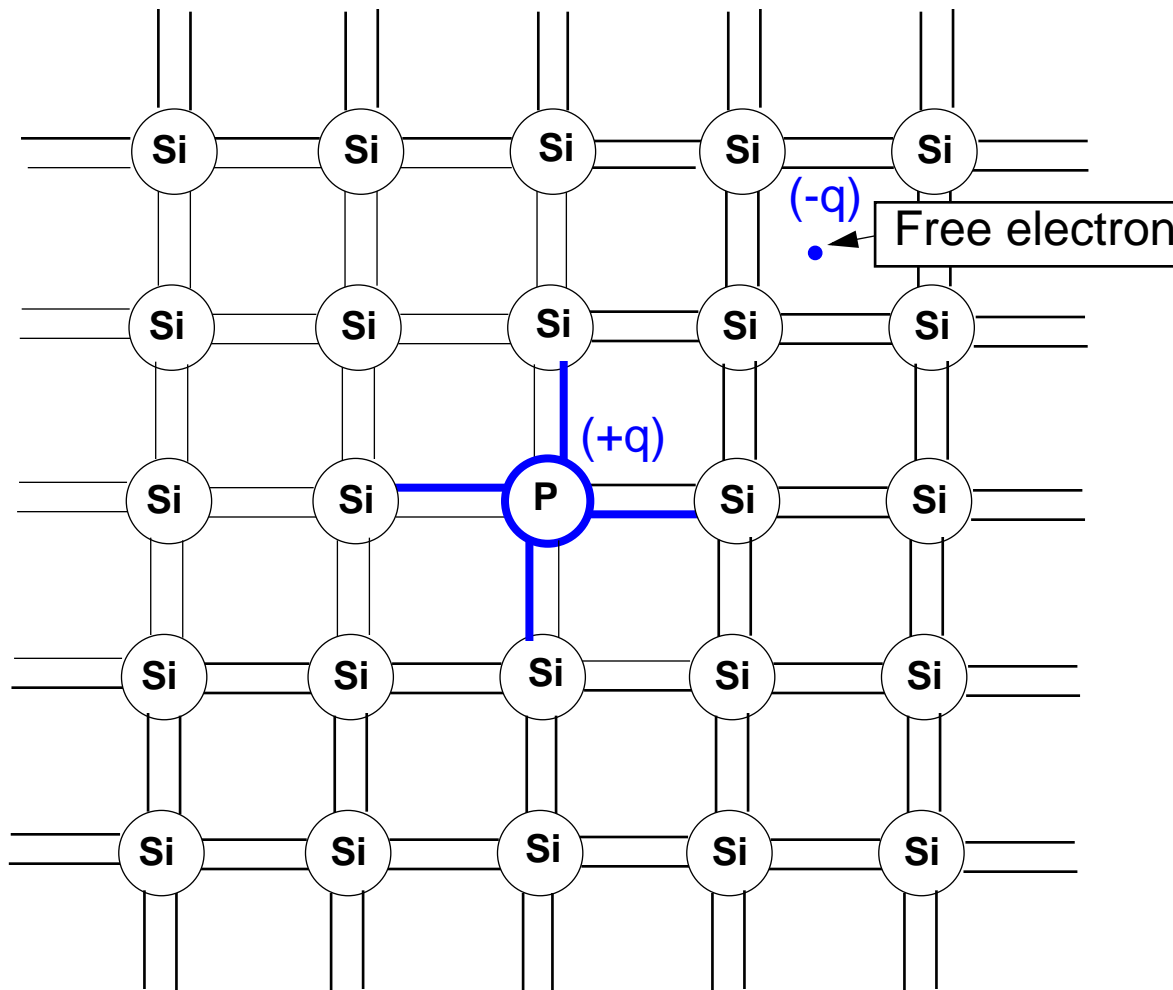
$$J = q(\mu_n n + \mu_p p)E$$



- What happens with the resistance when we increase carrier concentration?

## Doped Silicon(1) - Donor

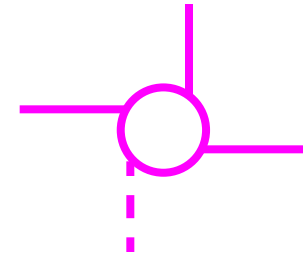
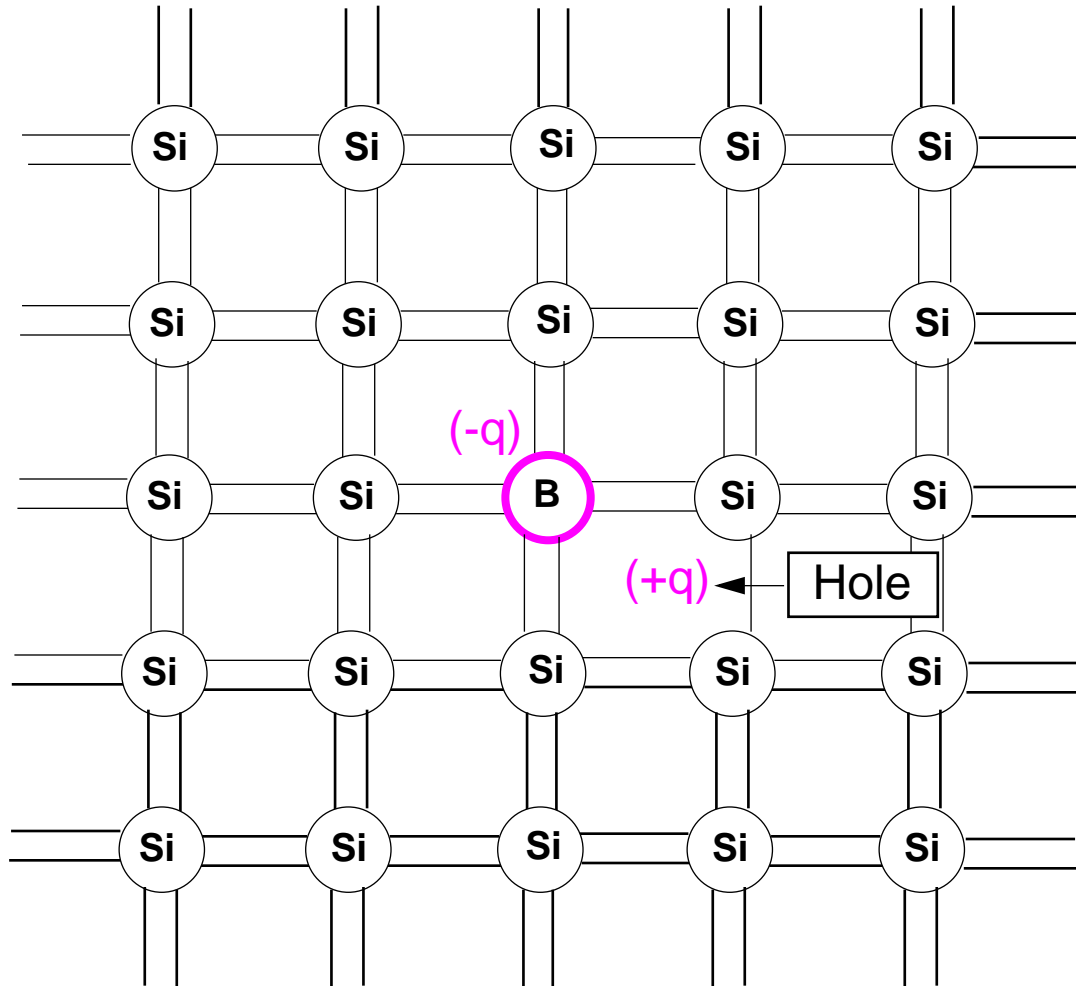
Donor has “one electron too much” - four of its electrons are bound. The excess electron is “donated” as a free electron. The donor is ionized.



- There is an increase in free electrons which will tend to recombine with the holes. There are more electrons than holes in **n type** Si.
- **Donor** ions are immobile.
- Free electrons move.

## Doped Silicon(2) - Acceptor

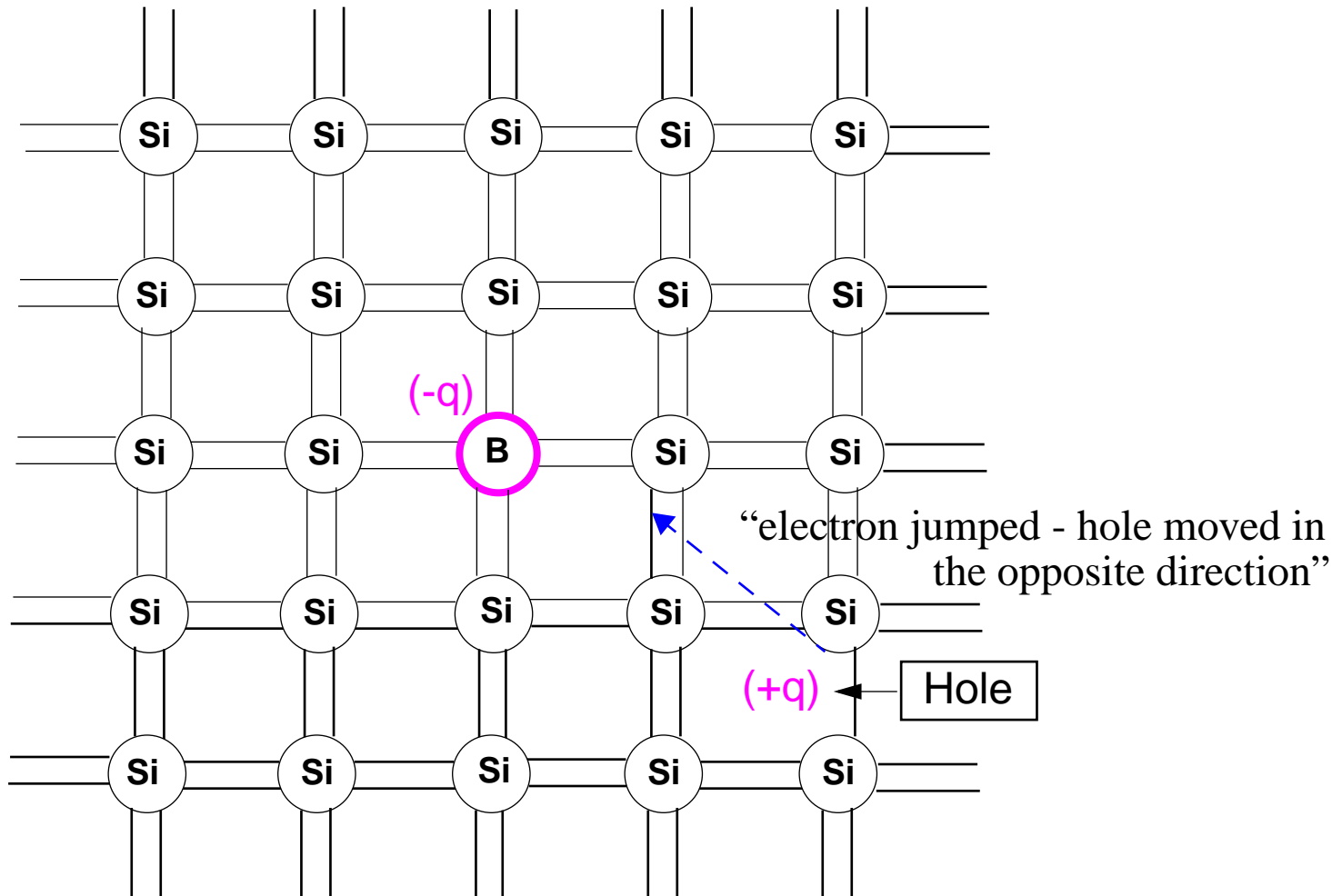
Donor has “one electron less than needed” - three of its electrons are bound. For the fourth bound, the donor catches a free electron. The donor is ionized.



- There is an increase in hole concentration since free electrons tend to “get trapped” by acceptors and leave the holes behind. There are more holes than electrons in **p type Si**.
- **Acceptor** ions are immobile.
- Holes move.

## Doped Silicon(3) - Holes Move

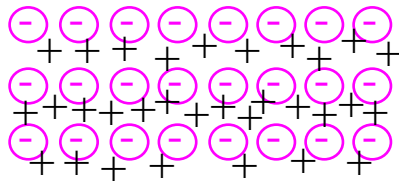
An electron may “jump” into the place where the hole is, but this will leave a hole somewhere else.



## Doped Silicon (4) - p-type, n-type

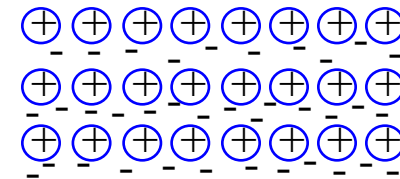
You can think of a doped silicon as of a material which:

- has a lot of “charged nodes” (impurities, dopants) built-in into the crystal lattice,
- is filled with gas of free charge carriers



p-type:

- vast majority of charge carriers are holes,
- negative built-in charge from acceptors.



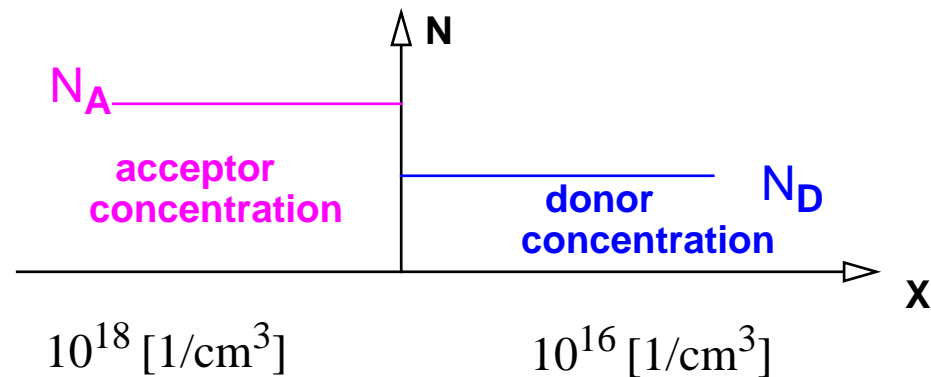
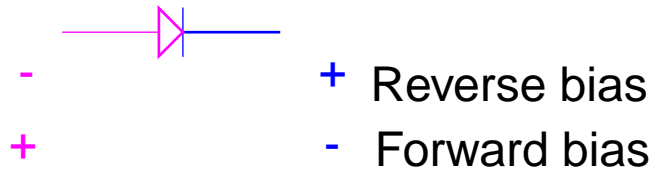
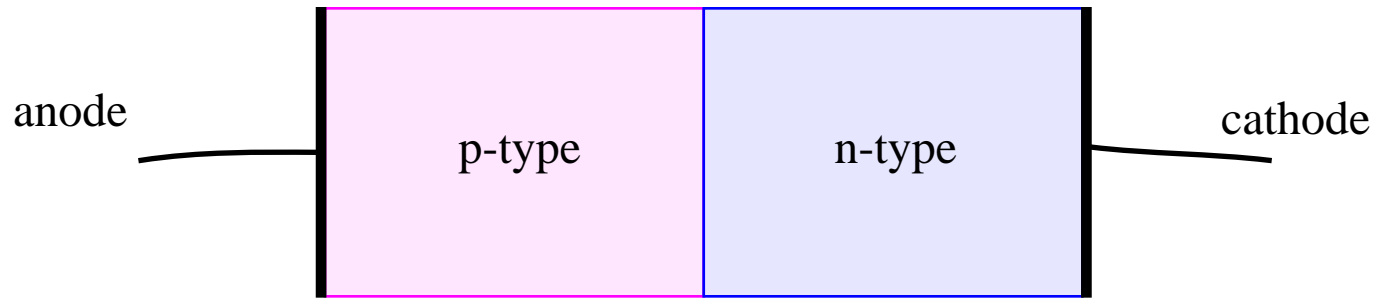
n-type:

- vast majority of charge carriers are electrons,
- positive built-in charge from donors.

- Note that even with doping, the impure silicon is still **charged neutral** since:
  - every “donated” free electron leaves behind a **positively charged donor** atom,
  - every **negatively charged acceptor** atom caused creation of a hole.

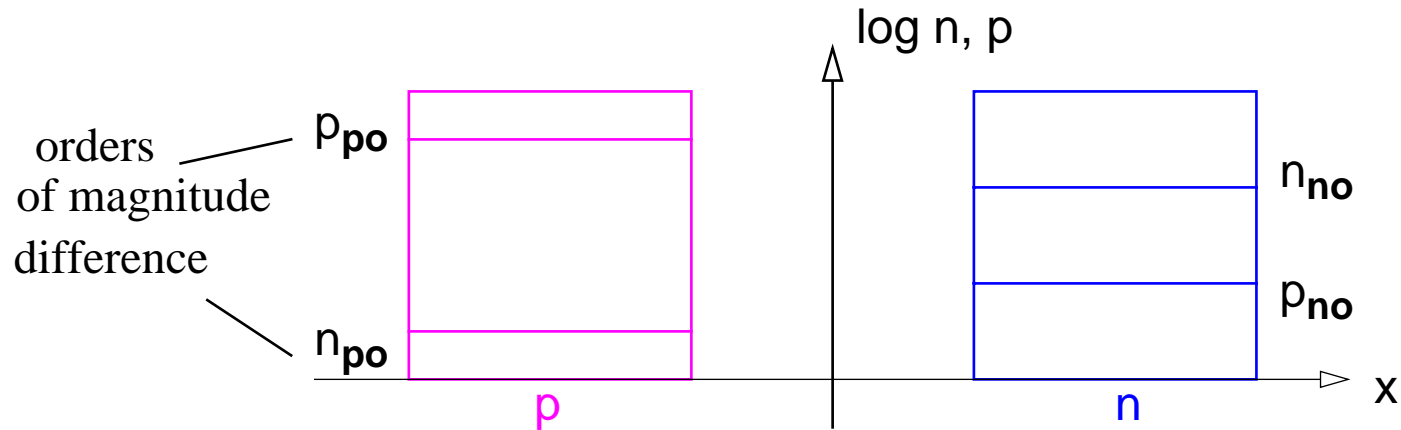
# p-n Junction Diode

Generally formed by introducing p-type impurities into n-type silicon, or vice versa



# p-n Junction - Example carrier concentrations

In equilibrium:  $p_{po} n_{po} = n_i^2 = n_{no} p_{no}$



**p-type region:**

$$N_A^- - p_{po} + n_{po} = 0$$

$$p_{po} \gg n_{po} \Rightarrow p_{po} = N_A^-$$

$$n_{po} =$$

**n-type region:**

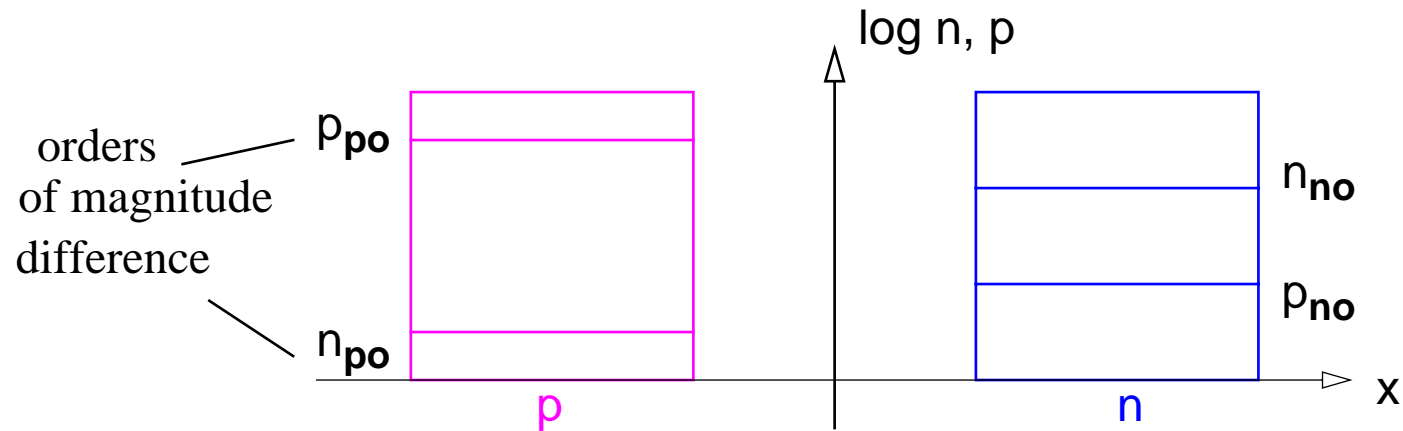
$$N_D^+ - n_{no} + p_{no} = 0$$

$$n_{no} \gg p_{no} \Rightarrow p_{no} = N_D^+$$

$$p_{no} =$$



# p-n Junction - Example carrier concentrations (2)



**p region** {

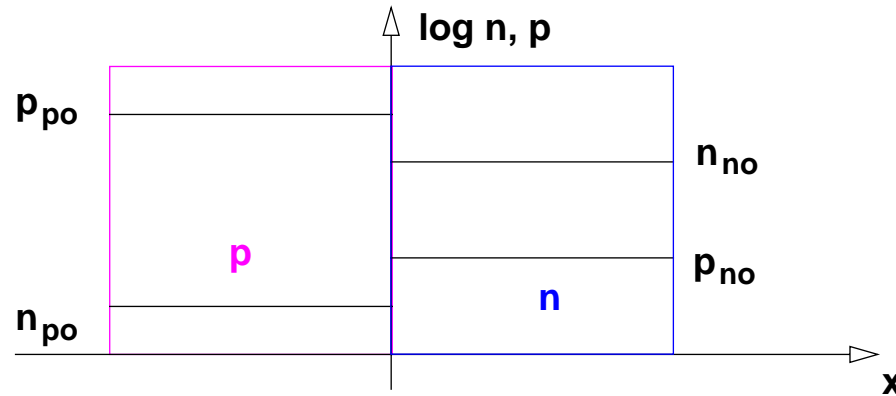
- majority carriers:  $p_{po} \approx N_A$   $10^{18}$
- minority carriers:  $n_{po} = \frac{n_i^2}{p_{po}} \approx \frac{n_i^2}{N_A}$  200

**n region** {

- majority carriers:  $n_{no} \approx N_D$   $10^{16}$
- minority carriers:  $p_{no} \approx \frac{n_i^2}{N_D}$   $10^4$

For comparison: how many Si atoms are in  $1 \text{ cm}^3$  ?

## Let us join the p and n regions

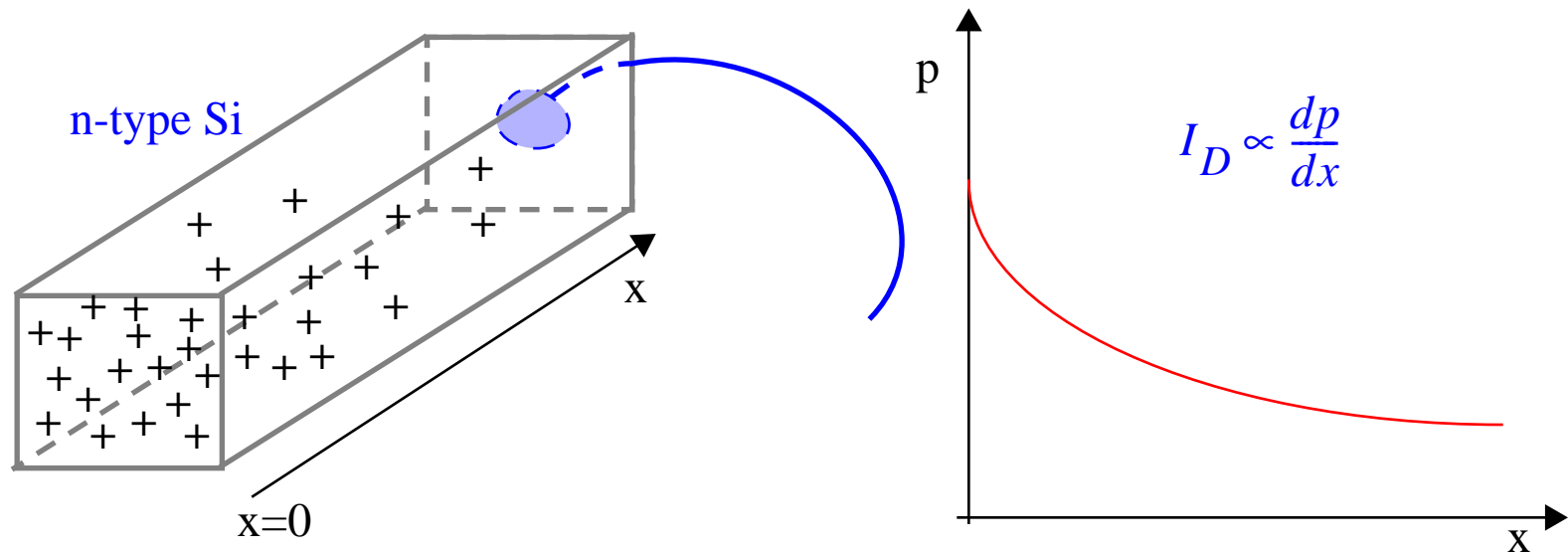


- At the interface enormous gradients of electron and hole concentrations.
- The holes will move towards n - region, electrons towards p - region. Acceptor and donor ions cannot move.
- Why will holes and electrons move to the opposite regions?

# Diffusion Current

- Holes enter the n-side then diffuse toward equilibrium
- If the source of holes were constant, this would continue indefinitely, with holes continually recombining with electrons

$$J_p \text{ (diff)} = -qD_p \frac{dp}{dx}$$



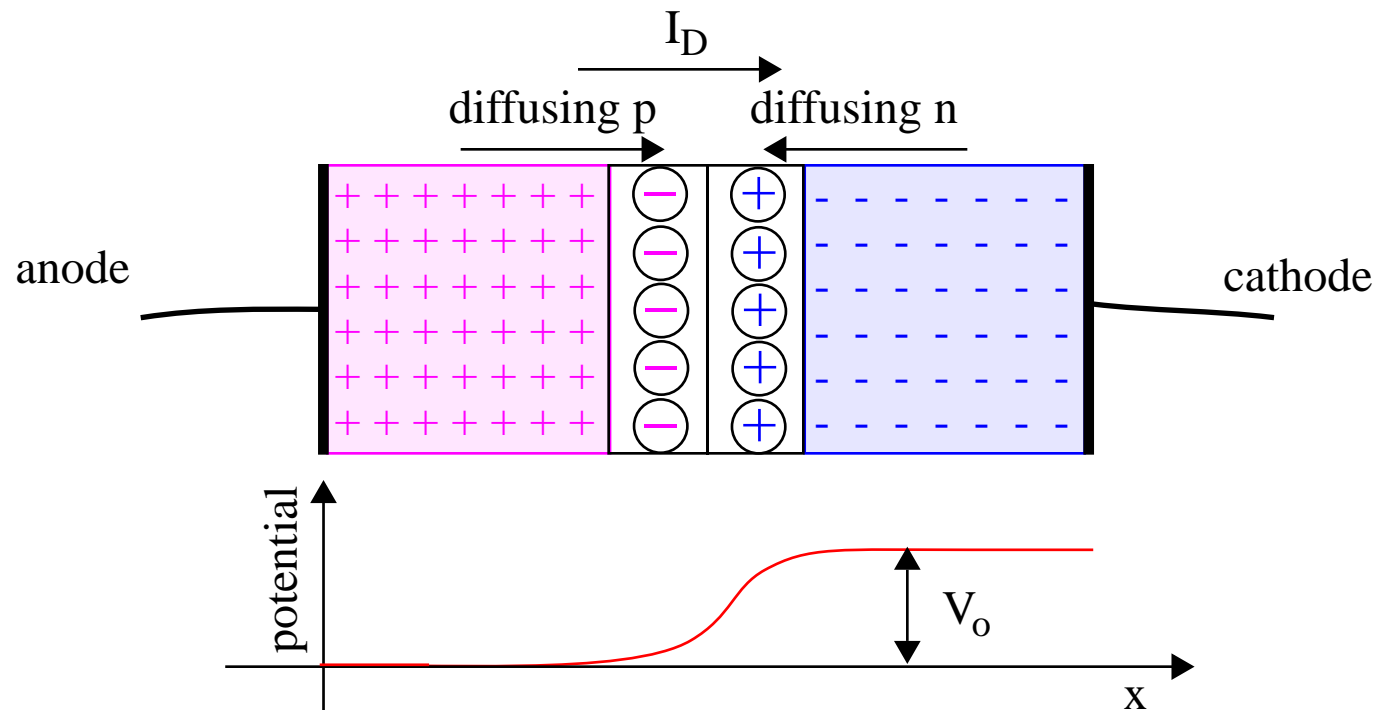
- There is a similar component of diffusion current due to electrons diffusing from the n-side toward the p-side

$$J_n \text{ (diff)} = qD_n \frac{dn}{dx}$$

- Diffusion currents do not result from electrostatic interactions!

## Diffusion Current and Depletion Region

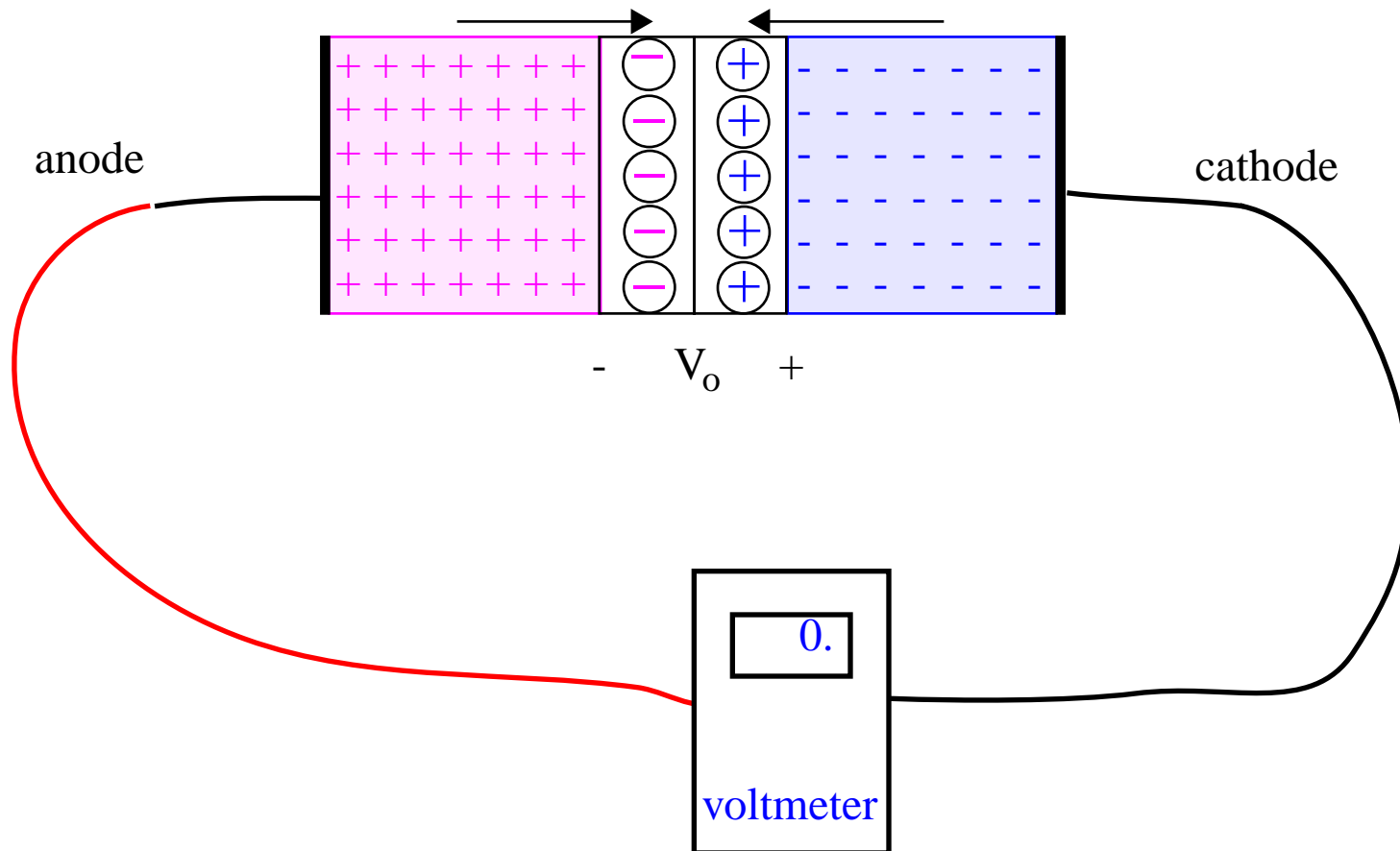
- Diffusing holes and electrons quickly recombine on the other side due to the high concentration of opposite free carriers there
- As they recombine they leave behind fixed charges due to uncovered bond charges on the donor and acceptor atoms
- This creates a **depletion region**



- The depletion region establishes an electric field
- Does this E-field oppose or aid the diffusion of holes and electrons?

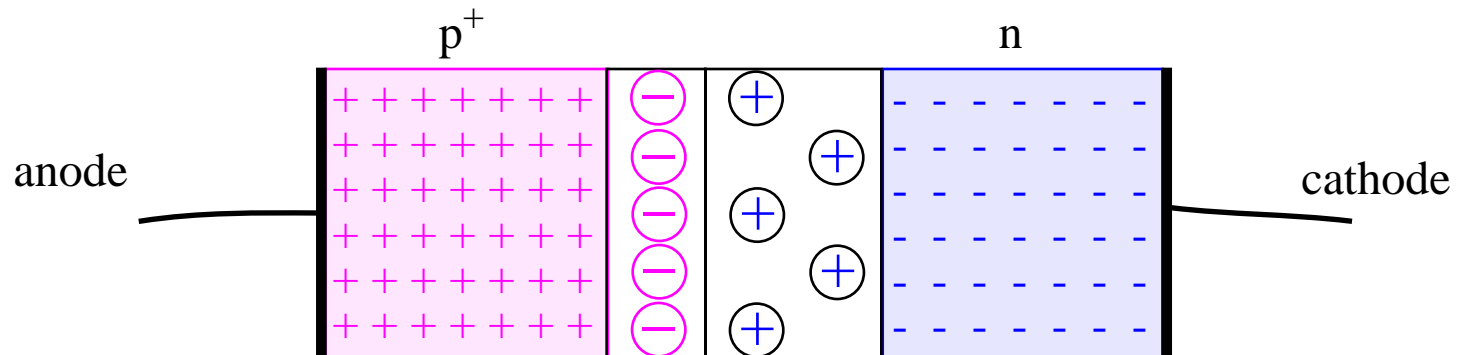
# Barrier Voltage, $V_0$

- The barrier voltage cannot be measured with a voltmeter
  - cannot draw energy from a p-n junction
- Connecting it to a meter would register 0 volts, which would indicate that the  $V_0$  voltage is dropped across the metal-Si contact points

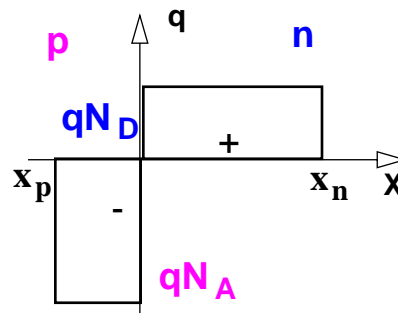


# Depletion Region

- One of the concentrations, p or n, is generally much larger than the other ( $p^+$  or  $n^+$ )



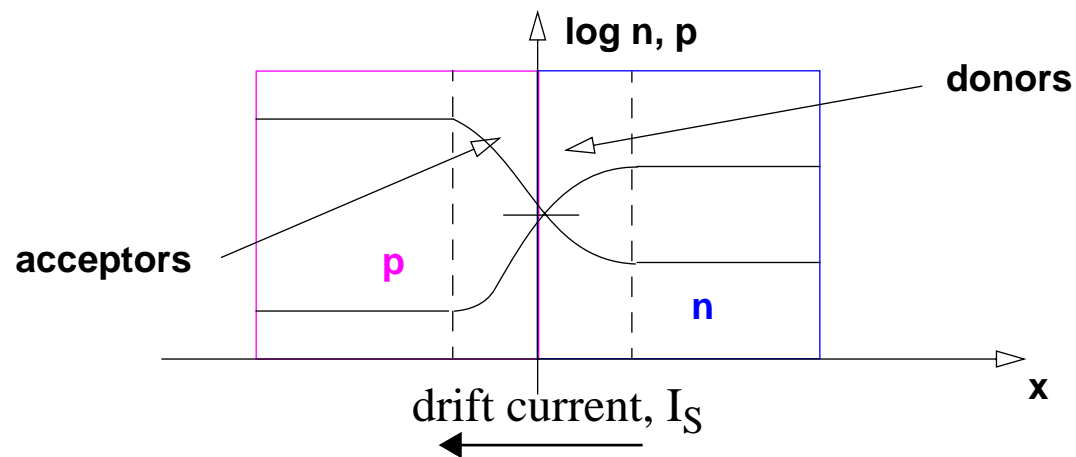
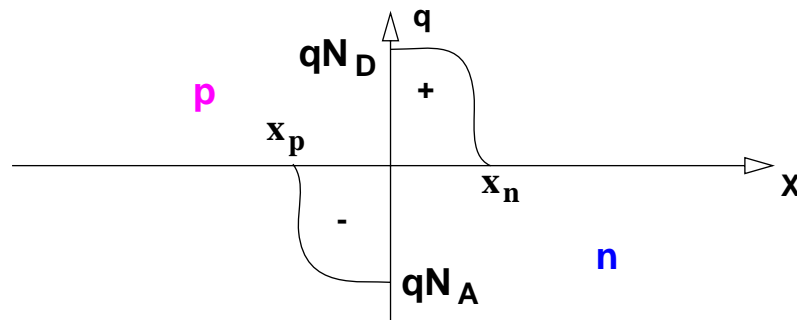
(Simplified plot for uniform doping)



- The depletion region extends almost entirely into a lighter doped side of the junction.

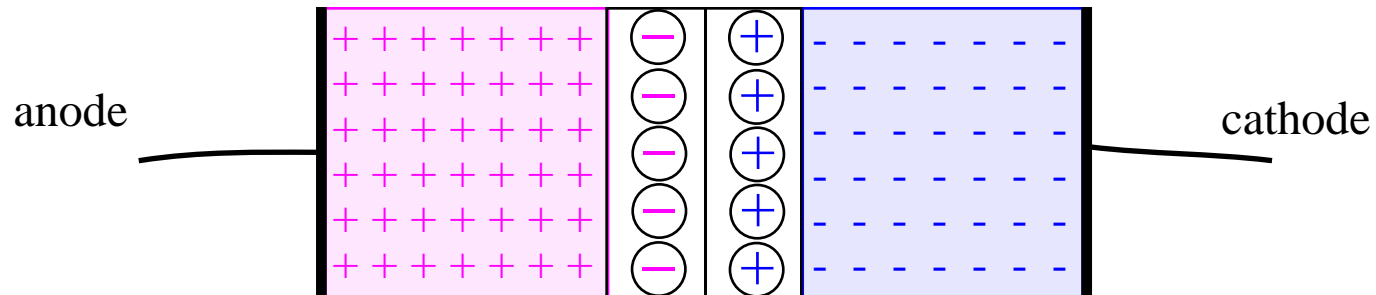
# Drift Current

- The electric field in the depletion layer forces electrons to flow toward the n-region and holes to flow toward the p-region. Thus, diffusion and drift currents flow in opposite directions.
- The few minority carriers (holes on the n-side and electrons on the p-side) that wander to the depletion region are quickly swept to the other side due to the electric field.



# Equilibrium

diffusing n  
drifting n  
diffusing p  
drifting p



$$I_S = I_p (\text{drift}) + I_n (\text{drift})$$

←

$$I_D = I_p (\text{diff}) + I_n (\text{diff})$$

→

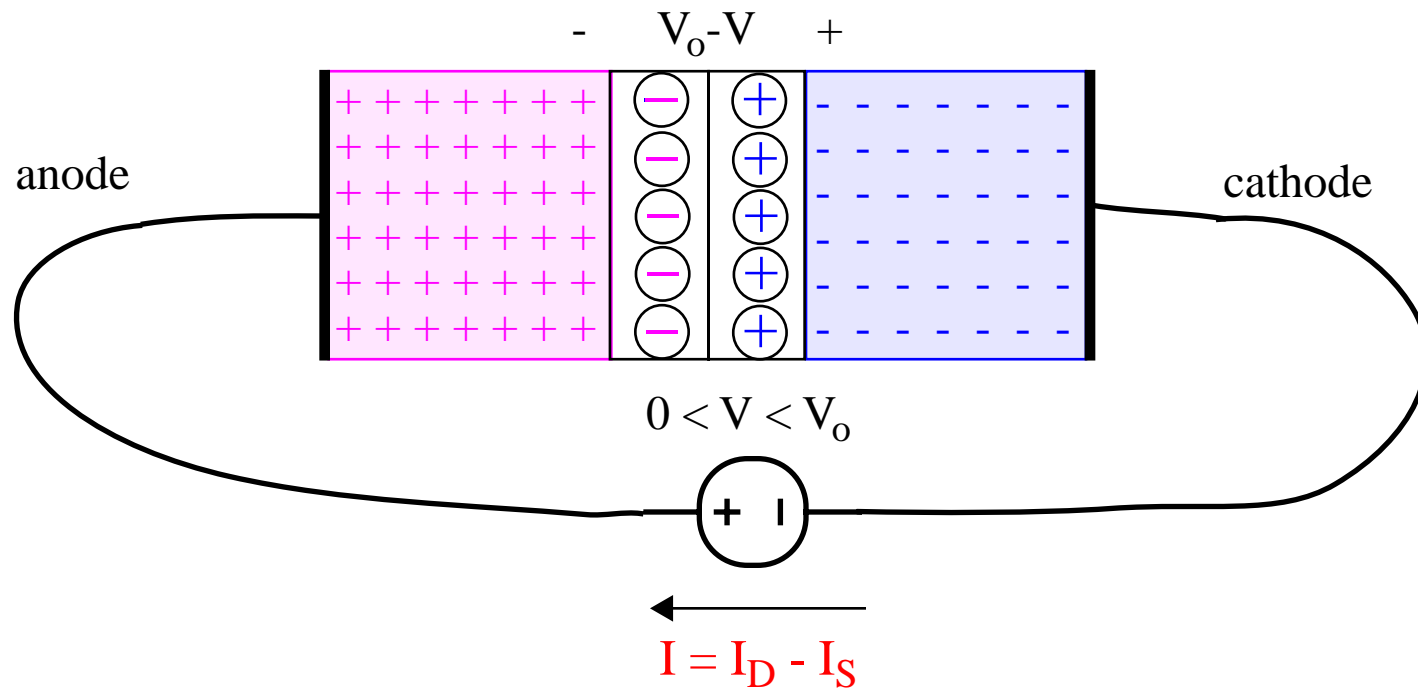
At equilibrium, these currents cancel each other and there is no net current flow:

$$I_S + I_D = 0 = I_p (\text{drift}) + I_n (\text{drift}) + I_p (\text{diff}) + I_n (\text{diff})$$



## Positive Applied Voltage

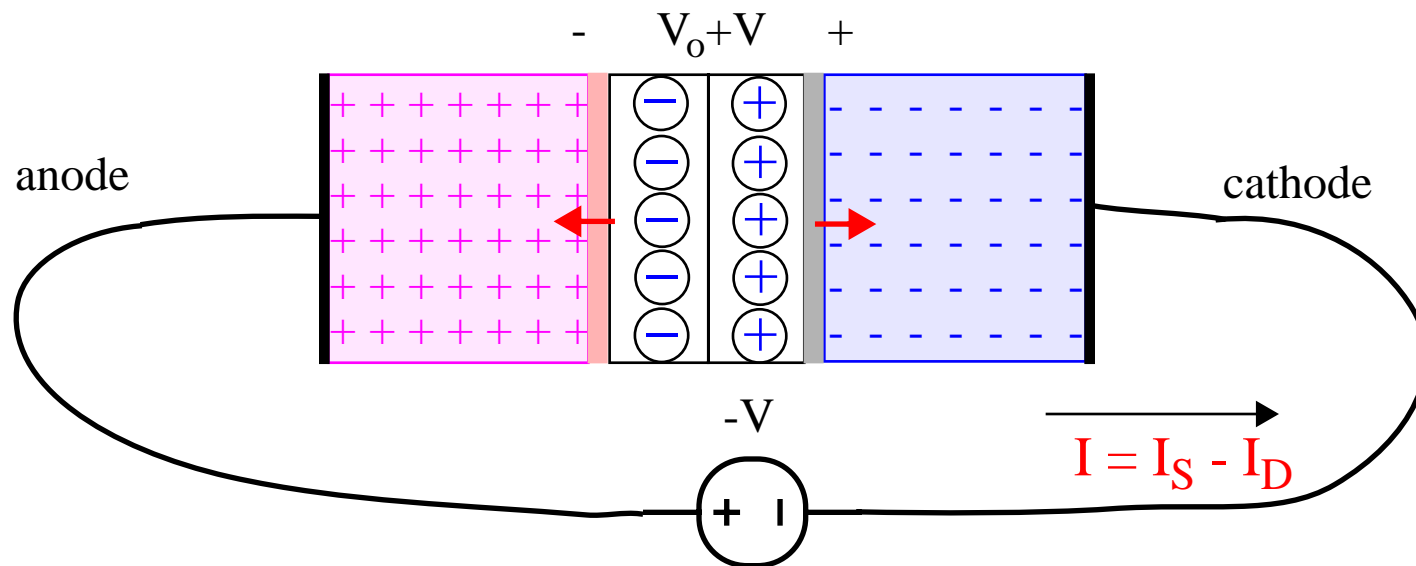
- A positive external voltage will reduce the barrier and allow more carriers to diffuse. Current can be large, because the carriers flow from the regions where their concentrations are large.
- The depletion region width is also reduced



- Now the diffusion current over the depletion region dominates the drift current

## Negative Applied Voltage

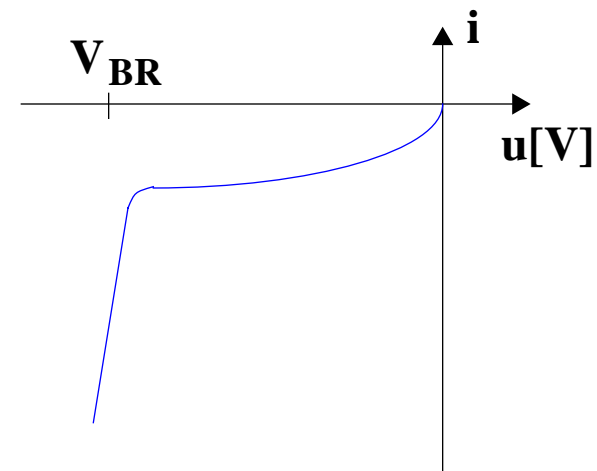
- A negative applied voltage will increase the width of the depletion region and increase the drift component of current,  $I_S$ . The diffusion current is decreased- Why?



- This current is small (e.g.  $10^{-12}$  A), because the carriers flow from the regions in which their concentrations are low.
- $I_S$  often negligible, depending on the circuit
- But there is a change in stored charge --- capacitance

# Breakdown Voltage

- **Zener breakdown:** negative applied voltage is so large that E-field is huge enough to break atom bonds and create lots of electron/hole pairs
- Electric field then sweeps these carriers to appropriate sides of diode thereby creating a large current
- Voltage remains constant at about  $V_Z$
- Diodes used for voltage regulation are designed for achieving a particular breakdown voltage
- **Avalanche breakdown:** E-field is so strong that electrons reach velocities that are fast enough that their collisions with atoms create new electron-hole pairs, which create more electron-hole pairs, and so on...



## Symbols used in this lecture

$N_D$	- donor concentration	$[cm^{-3}]$
$N_A$	- acceptor concentration	$[cm^{-3}]$
$J_p$	- hole current density	$[A/cm^2]$
$J_n$	- electron current density	$[A/cm^2]$
$q$	- electron charge (magnitude)	$[1.6 \times 10^{-19} \text{ coulombs}]$
$\mu_p$	- hole mobility	$[cm^2/volt \cdot s]$
$\mu_n$	- electron mobility	$[cm^2/volt \cdot s]$
$p$	- concentration of holes	$[cm^{-3}]$
$n$	- concentration of electrons	$[cm^{-3}]$
$D_p$	-diffusion coefficient for holes	$[cm^2/s]$
$D_n$	- diffusion coefficient for electrons	$[cm^2/s]$
$\frac{dp}{dx}$	- gradient of hole concentration	$[cm^{-4}/s]$
$\frac{dn}{dx}$	-gradient of electron concentration	$[cm^{-4}/s]$
$p_{po}$	- majority hole concentration	$[cm^{-3}]$

$n_{no}$	-majority electron concentration	$[cm^{-3}]$
$p_{no}$	-minority hole concentration	$[cm^{-3}]$
$n_{po}$	- minority electron concentration	$[cm^{-3}]$
$n_i$	- intrinsic concentration	$[1.4 \times 10^{10} cm^{-3}]$
A	- diode area	$[cm^2]$
$x_n$	- depletion region width in n-type semiconductor	$[cm]$
$x_p$	- depletion region width in p-type semiconductor	$[cm]$
$V_o$	- potential barrier	$[volt]$
$V_T$	- thermal voltage (26mV at room temperature)	$[volt]$
$k$	- Boltzmann's constant	$[1.38 \times 10^{-23} J/K]$
$V$	- bias voltage	$[volt]$
$T$	- absolute temperature $[273 + T(^{\circ}C)]$	$[K]$